

Direct ν mass measurements via beta-decay

Direct measurements of neutrino mass

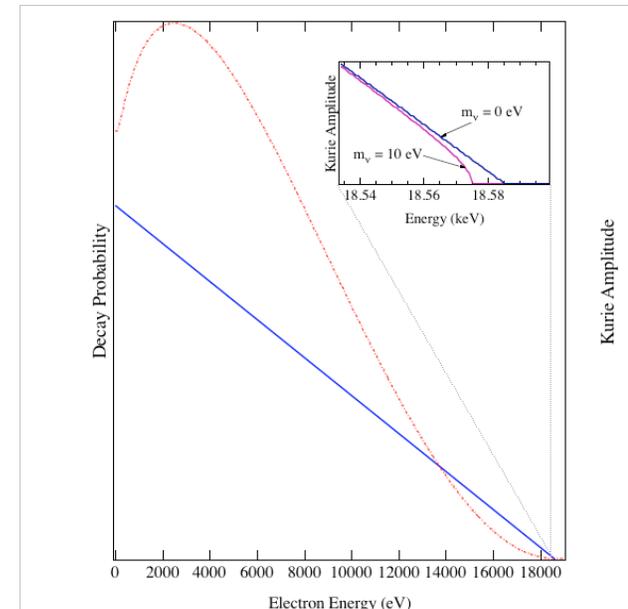
- Probe of absolute mass
- ν oscillations - a major paradigm shift
- ν -decay endpoint and ν mass eigenstates

Current ν -decay endpoint results

- INR - Troitsk
- Mainz

Future prospects for sub-eV sensitivities

- ^{187}Re bolometers
- Karlsruhe Tritium Neutrino Experiment (**KATRIN**)



Probes of absolute ν mass

- Indirect

- Cosmology - Hot Dark Matter

- Galaxy clusters
- Lyman- α forest
- Galaxy large scale structure

- Astrophysics

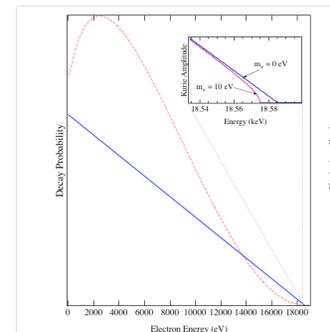
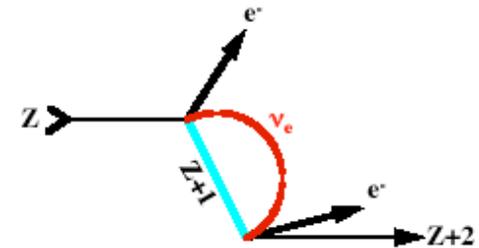
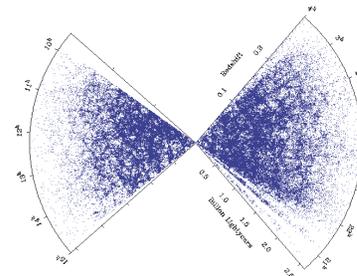
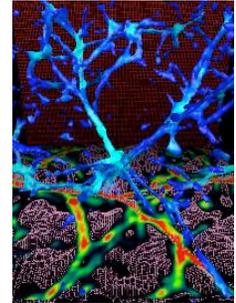
- UHE cosmic-rays
- Supernovae generation mechanisms

- Neutrinoless $\beta\beta$ -decay

- Direct techniques

- time of flight (supernovae)

- **particle decay kinematics**

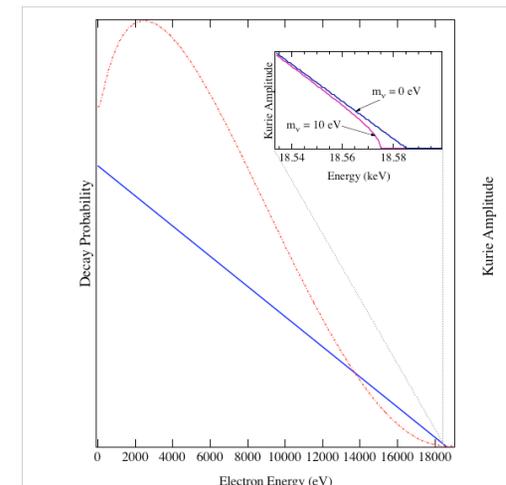
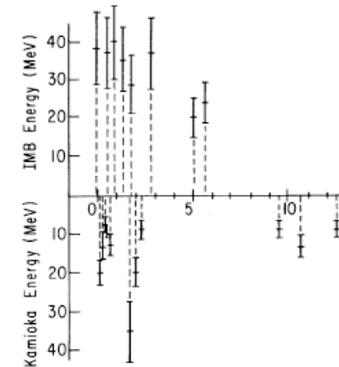


Direct measurements of neutrino mass

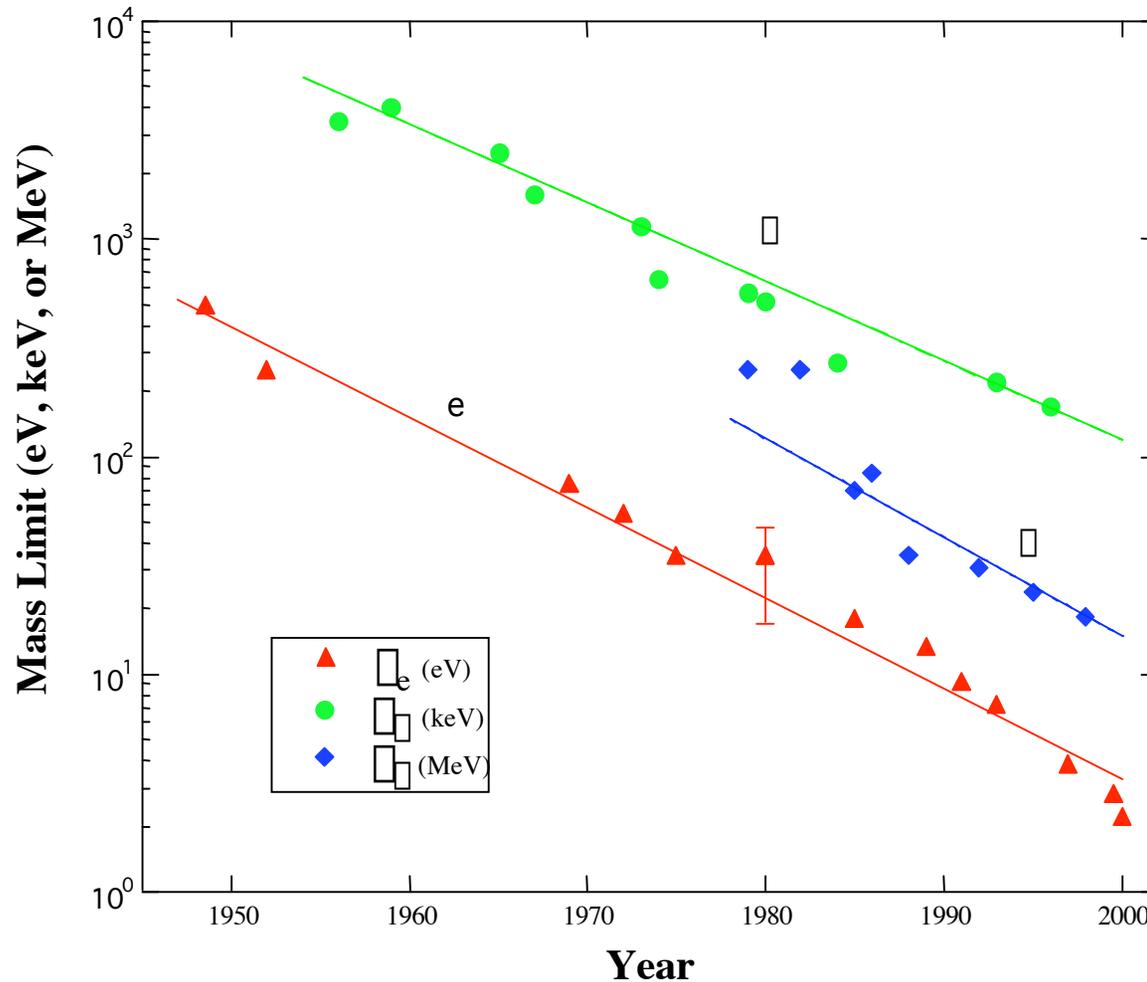
Techniques

- **time of flight from supernovae**
 - not expected to reach sub-eV sensitivity
 - Collapse to Black Hole: $m_{\nu} \geq 1.8$ eV (Rotation neglected)
Beacom, Boyd, Mezzacappa PRD63 073011
 - much more likely that a direct measurement or limit will be used to help understand supernovae dynamics.
- **particle decay kinematics**
 - β -decay (and e capture) spectrum shape
 - muon momentum in pion decay
 - invariant mass studies of multiparticle semileptonic decays
 - advantages
 - purely kinematical observables
 - few, if any, assumptions about ν properties
 - free of model dependencies

SN1987a



Past: history of direct μ mass measurements (μ flavor eigenstates)



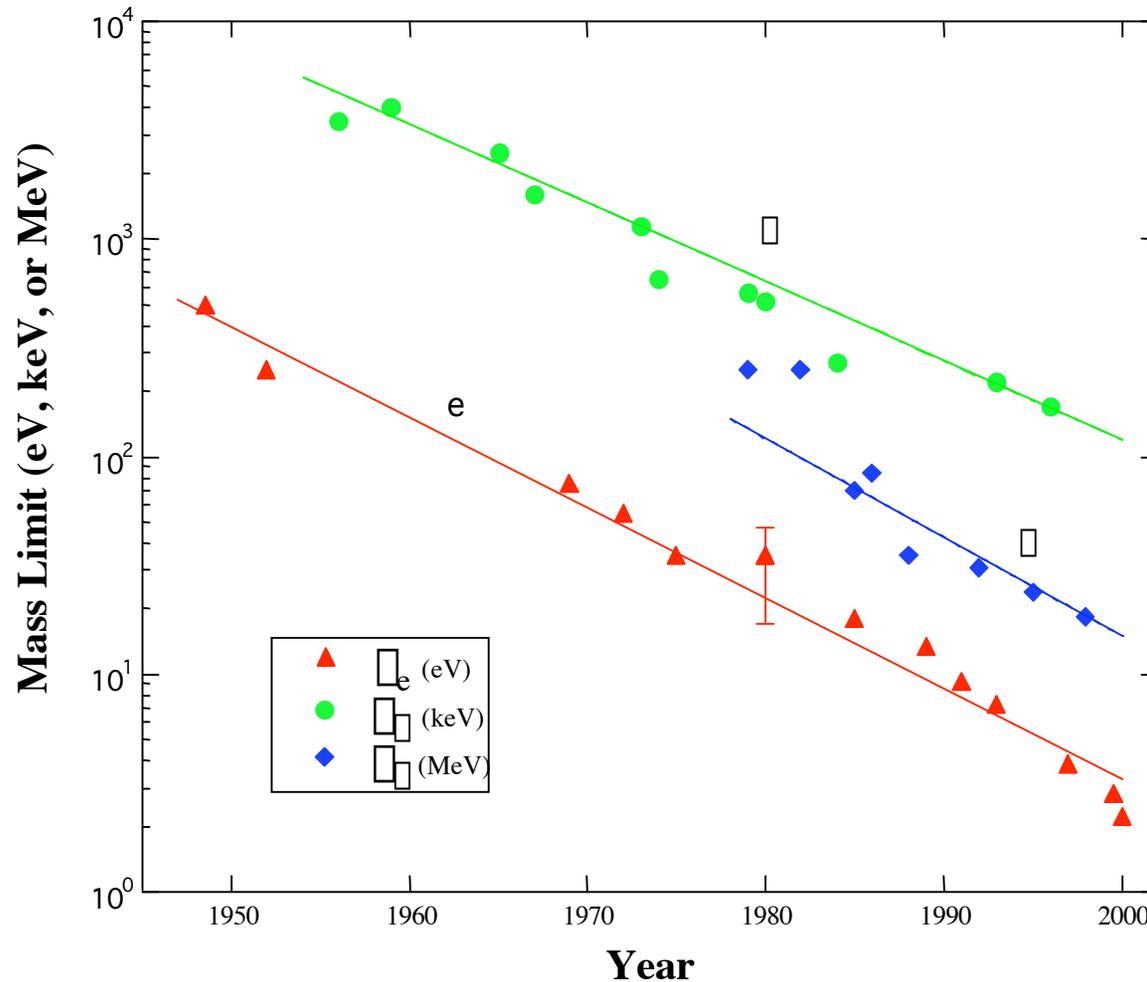
$m_{\mu\mu} < 170 \text{ keV}$ (90%CL)
(PSI 1996)

$m_{\mu\mu} < 18.2 \text{ MeV}$ (95% CL)
(ALEPH 1998)

$m_{\mu e} < 2.2 \text{ eV}$ (95% CL)
(Mainz 2000)

points without error bars represent upper limits

Past: history of direct ν mass measurements (ν flavor eigenstates)



But ν oscillations with large mixing angles - means one must consider direct techniques in terms of ν mass eigenstates!

points without error bars represent upper limits

□ mass & mixing - oscillation experiments

For a 3 neutrino scenario the lepton mixing matrix (Maki-Nakagawa-Sakata-Pontecorvo), which relates □ mass eigenstates to weak or flavor eigenstates, is:

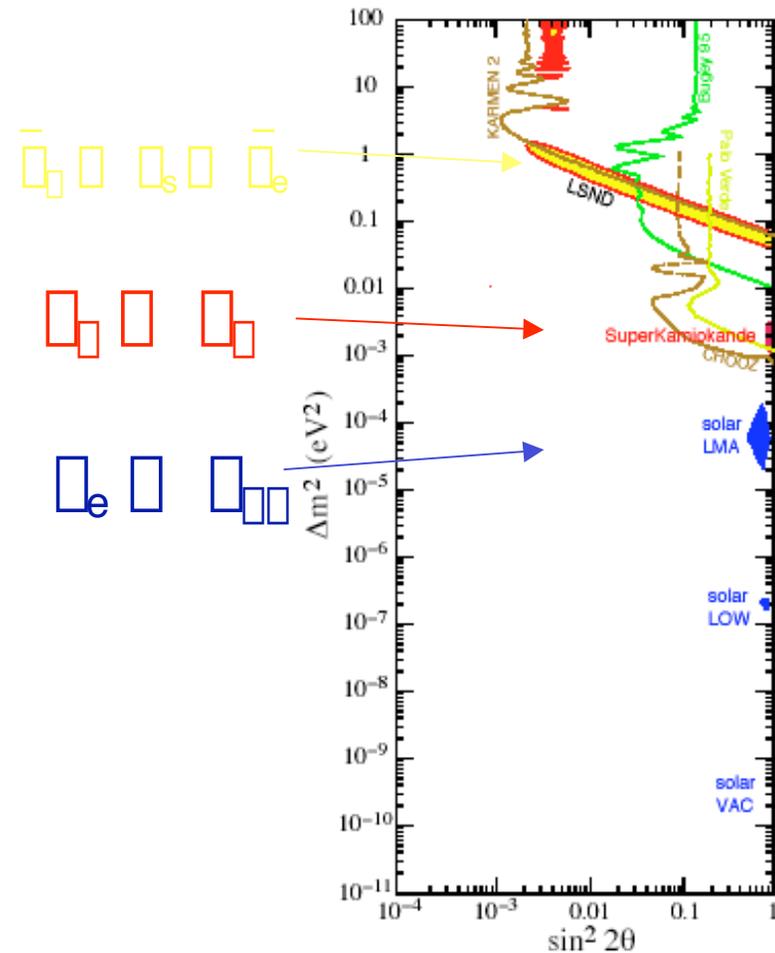
$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

Oscillation experiments yield mass squared differences:

$$\Delta m_{12}^2 = \left| m_{\nu_1}^2 - m_{\nu_2}^2 \right|$$

$$\Delta m_{23}^2 = \left| m_{\nu_2}^2 - m_{\nu_3}^2 \right|$$

$$\Delta m_{13}^2 = \left| m_{\nu_1}^2 - m_{\nu_3}^2 \right|$$



ν mass & mixing - oscillation experiments

Best bet for MNSP Matrix:

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

Atmospheric

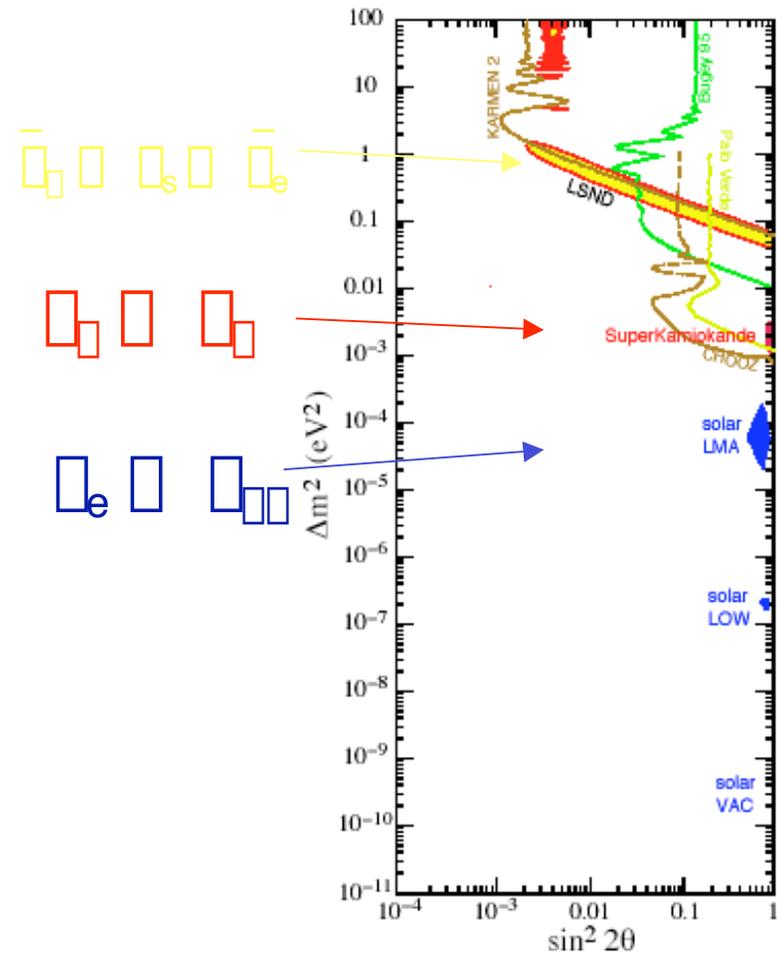
$$\begin{pmatrix} 1 & 0 & 0 \\ 0 & 1/\sqrt{2} & 1/\sqrt{2} \\ 0 & -1/\sqrt{2} & 1/\sqrt{2} \end{pmatrix} \times$$

Chooz

$$\times \begin{pmatrix} \sim 1 & 0 & e^{-i\delta_{CP}} \sin \theta_{13} \\ 0 & 1 & 0 \\ -e^{i\delta_{CP}} \sin \theta_{13} & 0 & \sim 1 \end{pmatrix} \times$$

LMA

$$\times \begin{pmatrix} 0.85 & 0.51 & 0 \\ -0.51 & 0.85 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$



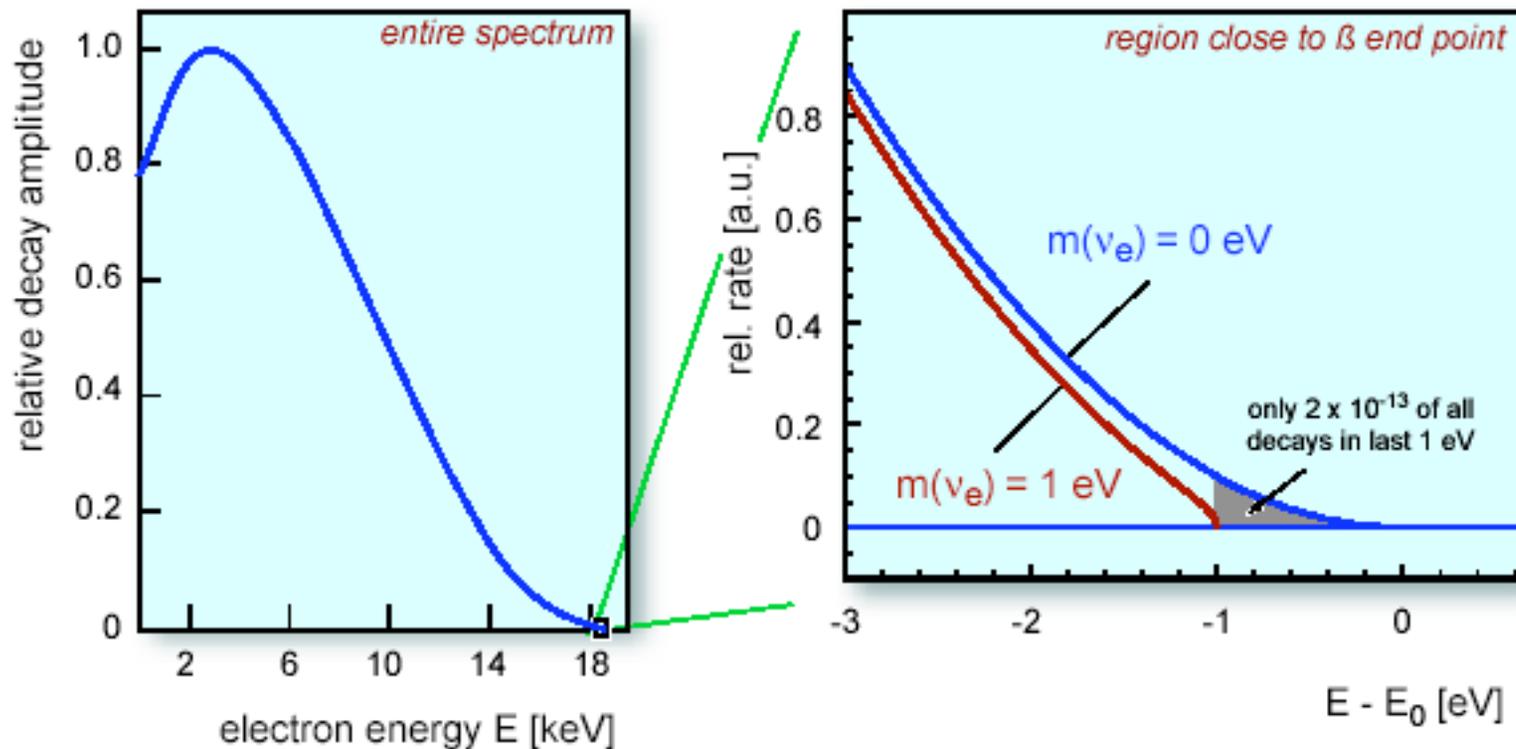
$$\nu_e = 0.85\nu_1 + 0.51\nu_2$$

$$\nu_\mu = 0.36\nu_1 + 0.60\nu_2 + 0.71\nu_3$$

$$\nu_\tau = 0.36\nu_1 - 0.60\nu_2 + 0.71\nu_3$$

β -decay endpoint measurement

Essentially a search for a distortion in the shape of the β -spectrum in the endpoint energy region



$$dN(E) = K |M|^2 F(Z, R, E) p_e E (E_0 - E) \left\{ (E_0 - E)^2 - m_{\nu_e}^2 c^4 \right\}^{1/2} dE$$

β -decay in terms of ν mass eigenstates

Taking into account ν mass eigenstates, the original spectrum

$$dN(E) = K|MI|^2 F(Z,R,E) p_e E (E_0 - E) \{(E_0 - E)^2 - m_{\nu_e}^2 c^4\}^{1/2} dE$$

becomes

$$dN(E) = K|MI|^2 F(Z,R,E) p_e E (E_0 - E) \sum_i |U_{ei}|^2 \{(E_0 - E)^2 - m_{\nu_i}^2 c^4\}^{1/2} dE$$

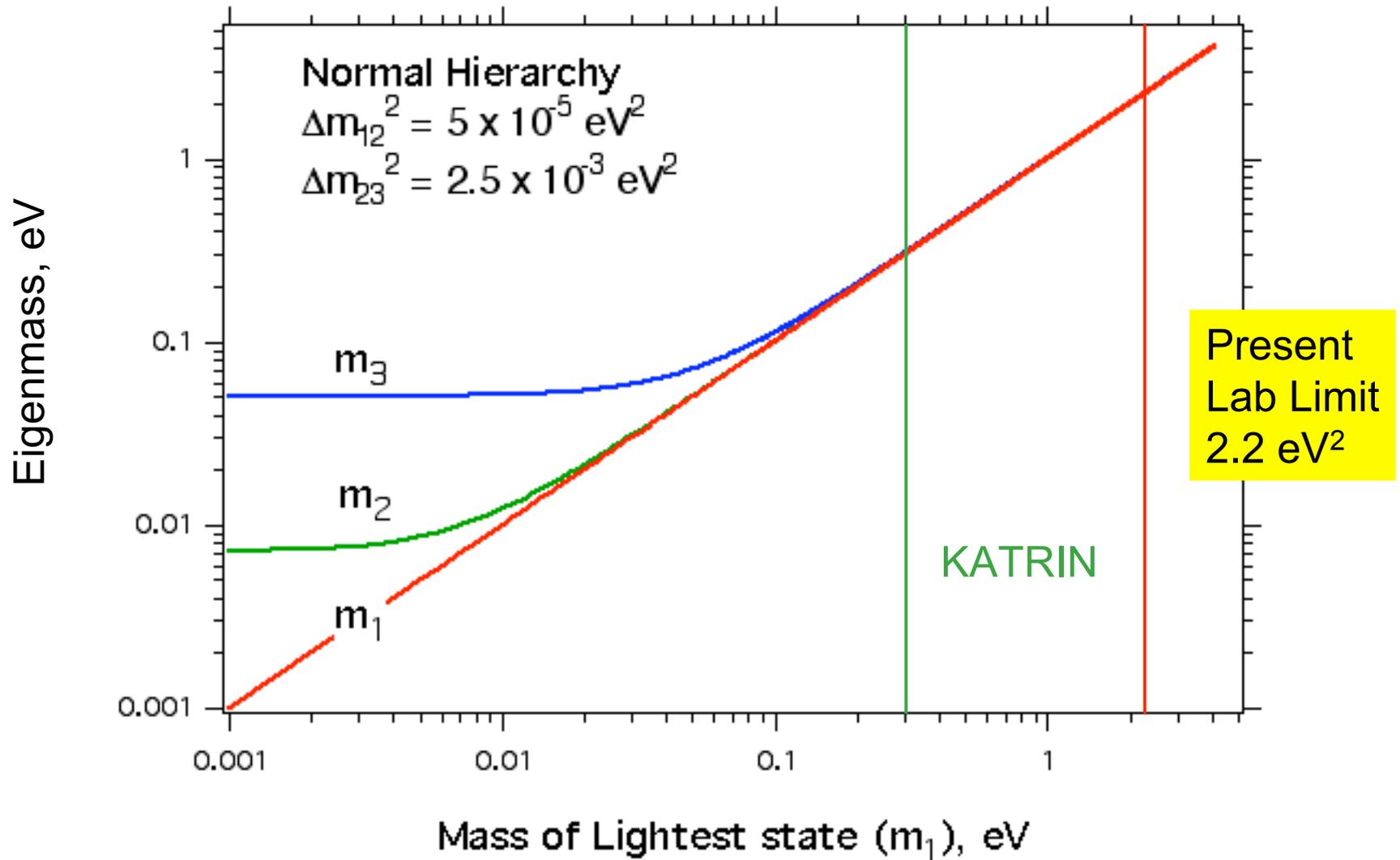
The observed beta spectrum shape will depend on:

- the neutrino masses
- the number of neutrino mass eigenstates
- the leptonic mixing matrix elements
- the total resolution/sensitivity of the measurement

For 3 ν mass spectrum, with degenerate states, the beta spectrum

simplifies to an “effective mass” : $m_{\nu}^2 = \sum |U_{ei}|^2 m_{\nu_i}^2$

All masses linked to lightest by oscillations



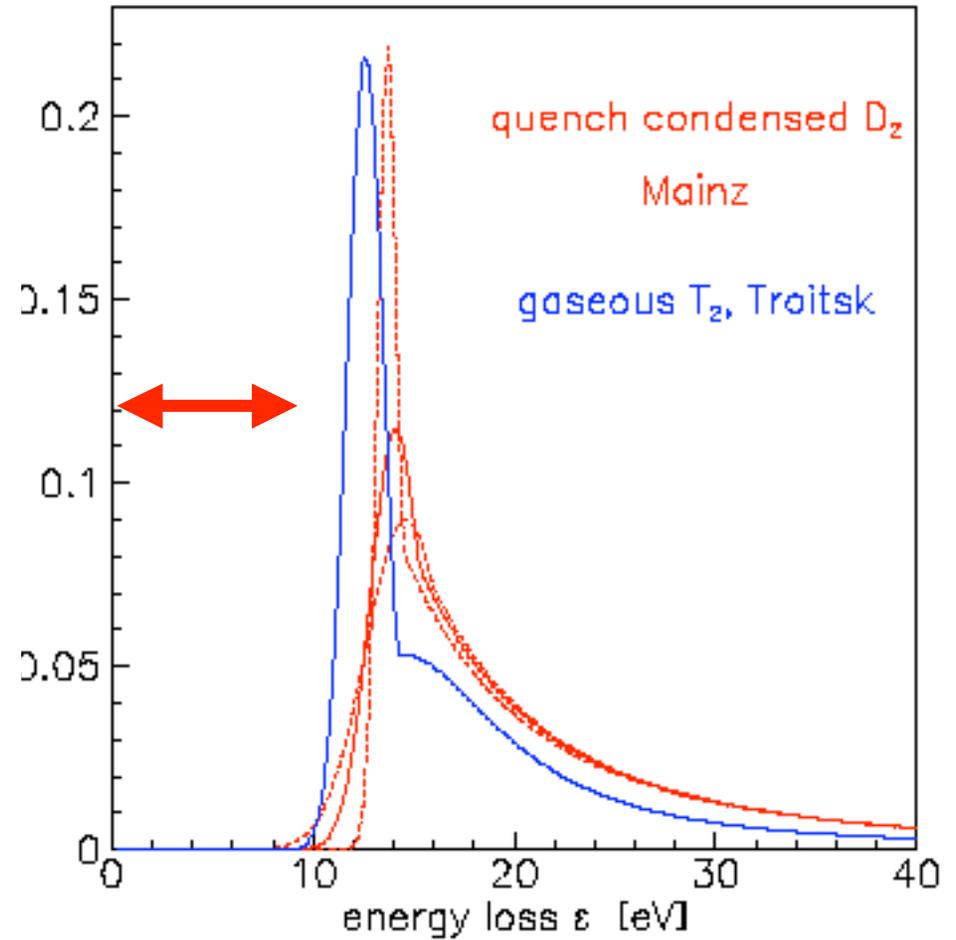
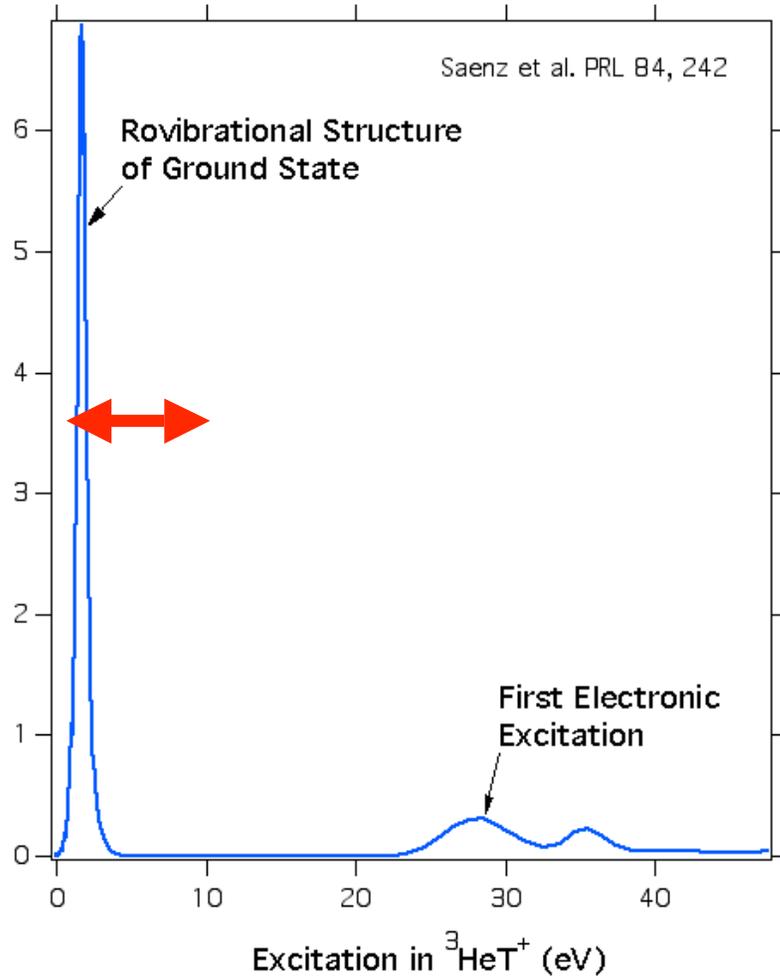
Keys to β -decay shape measurements

- **Statistics and uncertainty budget**
 - Only $2 \cdot 10^{-13}$ decays in last 1 eV below endpoint.
 - For 10 eV sensitivity, 100 eV², for 1 eV sensitivity, 1 eV²
 - Must reduce backgrounds (\sim mHz) and ensure that they are very stable with time.
- **One must precisely eliminate or characterize all possible shape effects**
 - atomic final state effects
 - use atomic or molecular tritium source (${}^3\text{H} \rightarrow {}^3\text{He} + e^- + \bar{\nu}_e$)
 - utilize spectrum above atomic states (last 20 eV below endpoint)
 - energy loss shape effects
 - directly measure
 - use only no-loss portion of spectrum (last 9 eV below endpoint)
 - instrumental shape effects
 - direct measurements, using ${}^{83}\text{Kr}^m$
 - use integral spectrometers with very good resolution (\sim eV)

A window to work in

Molecular Excitations

Energy loss function



Solenoid Retarding Spectrometer

Magnetic Adiabatic Collimation with Electrostatic Filter (MAC-E)

guiding by magnetic fields
(magnetic adiabatic collimation)

$$\Delta\Omega \sim 2\pi$$

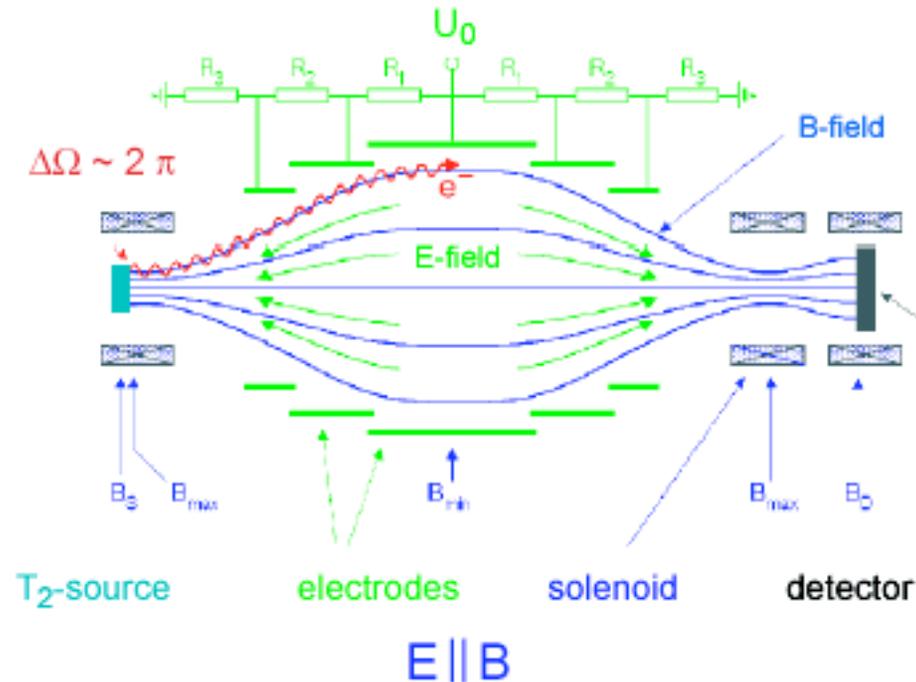
electric (retarding-) field :
analysis of electron energies
(electrostatic filter)

integral transmission : $E > U_0$

$$\vec{F} = (\vec{\mu} \cdot \vec{\nabla}) \vec{B} + q \vec{E}$$

$$\mu = E_{\perp} / B = \text{const}$$

adiabatic motion

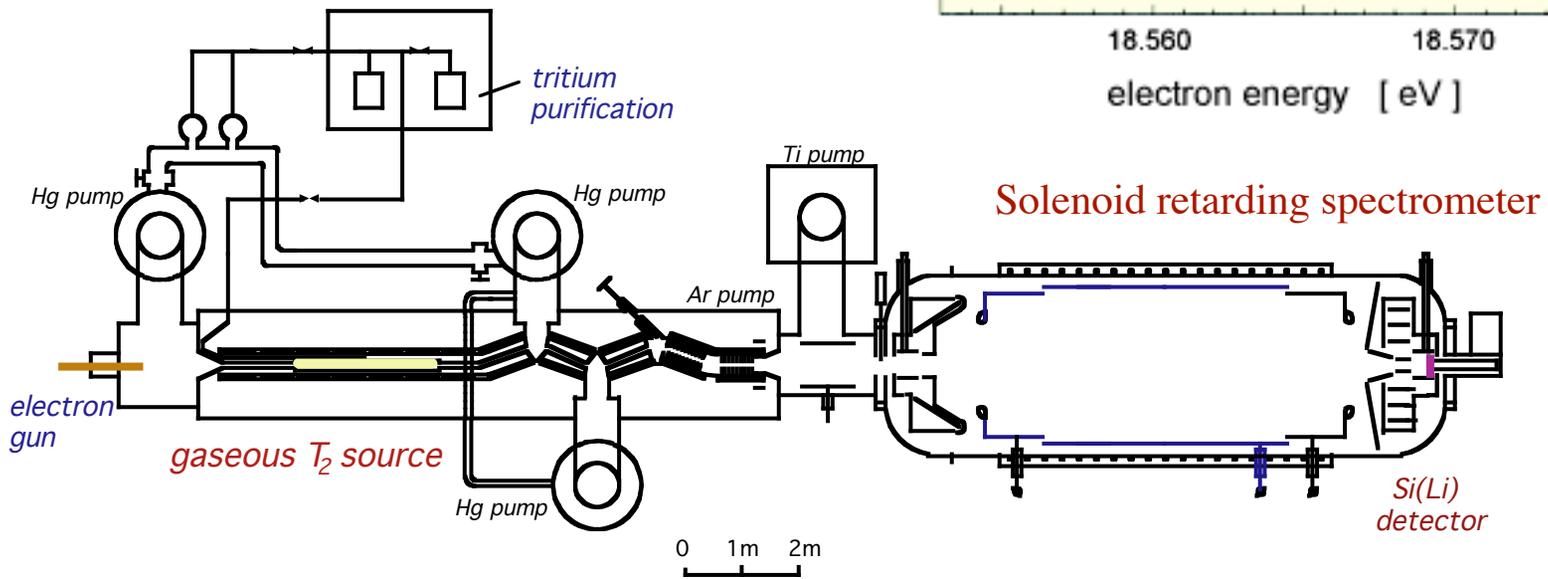
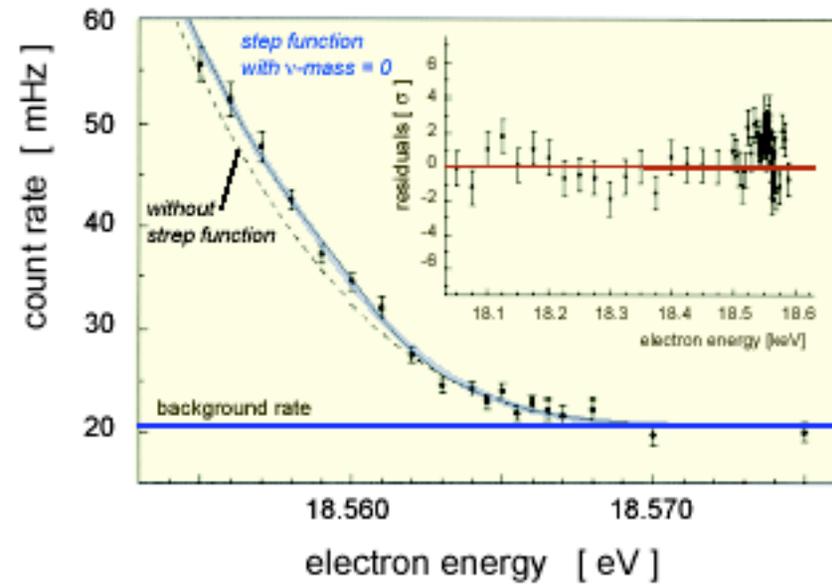


adiabatic transformation $E_{\perp} \rightarrow E_{\parallel}$

Troitsk tritium β -decay experiment



200 days of data since 1994



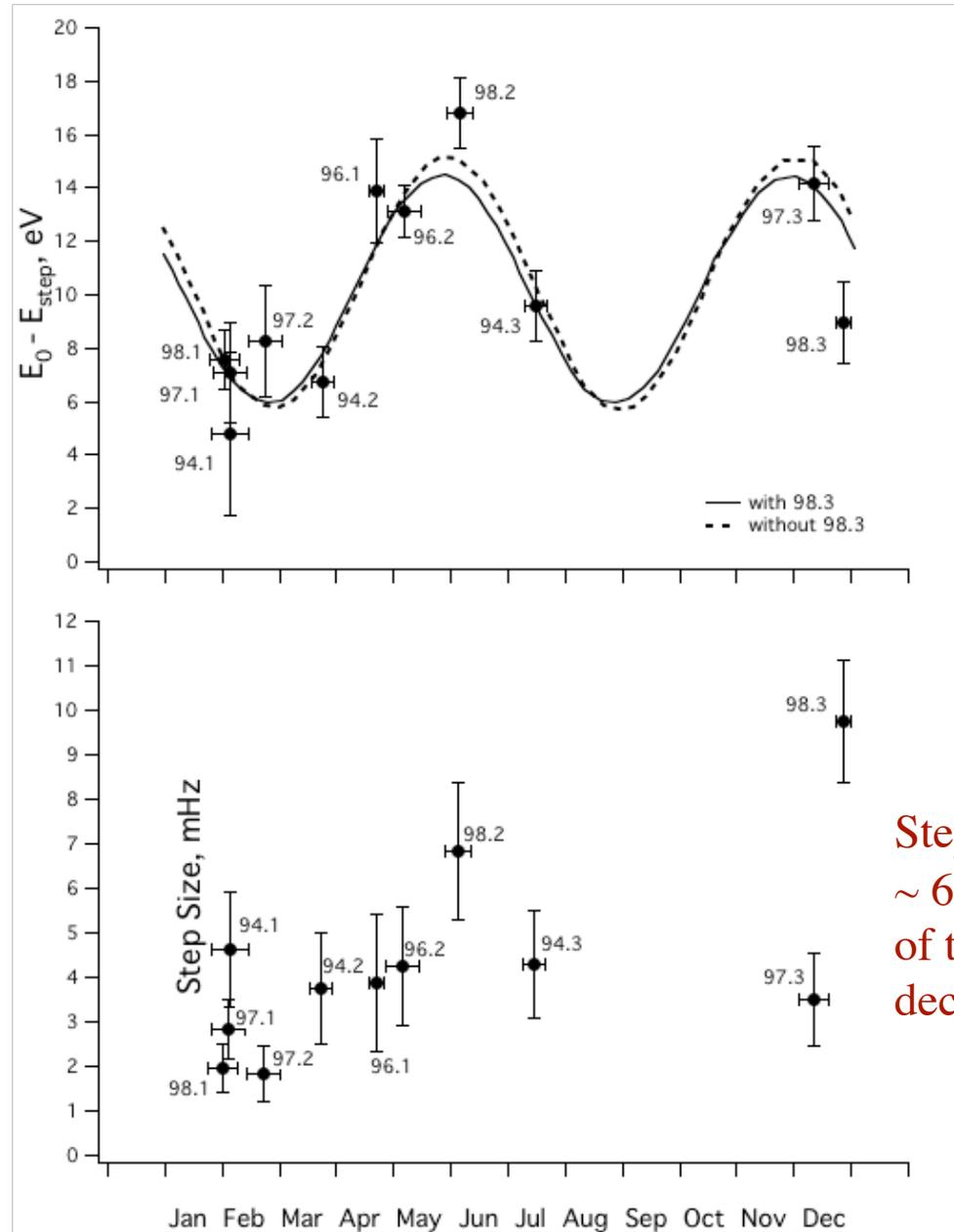
Troitsk Results

Claims there is a step function anomaly that varies in **both** amplitude and position above the endpoint.

It is difficult to have much confidence in their reported limit

$m_{\bar{\nu}} \leq 2.5 \text{ eV}$ (95%CL)
since it requires removing the step function (excess counts)

Likely systematic problems



Step intensity
 $\sim 6 \cdot 10^{-11}$
of total T_2
decay rate

Mainz Neutrino Mass Experiment

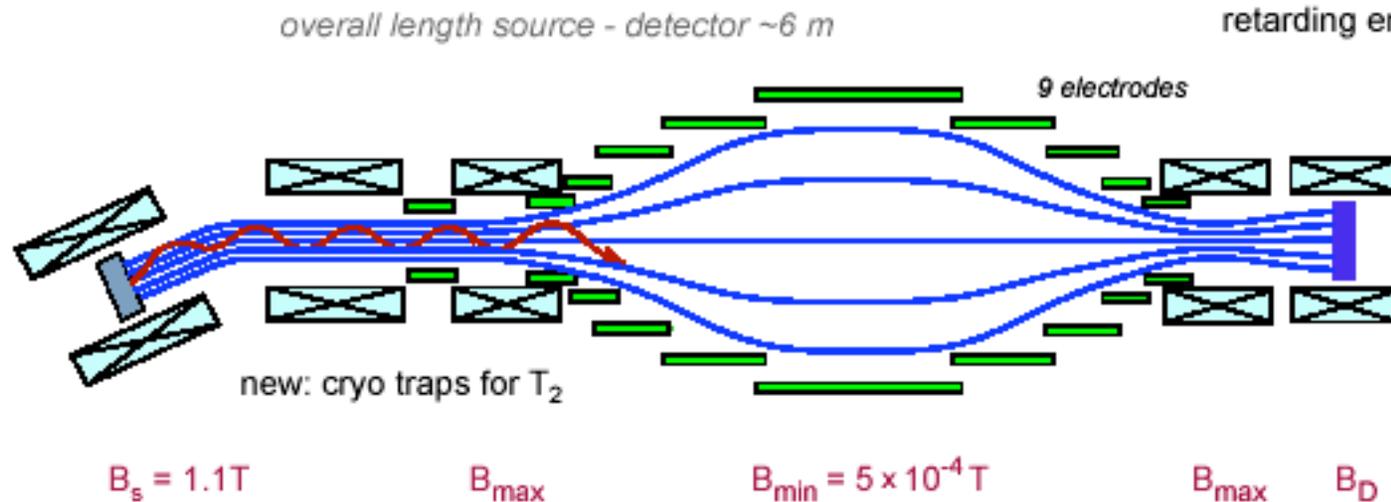
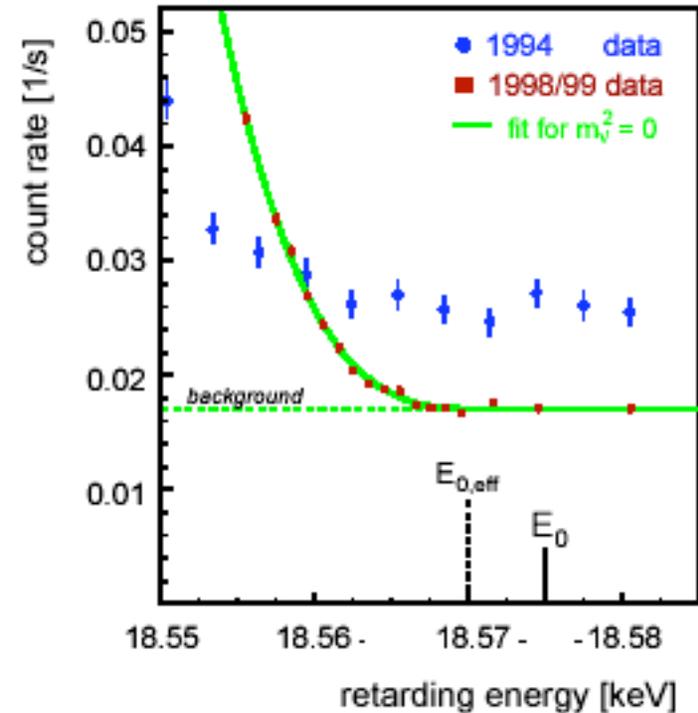


Quench condensed solid T_2 source

Early results (94) showed systematic effects, traced to source film roughening transition.

(fixed by lowering temperature)

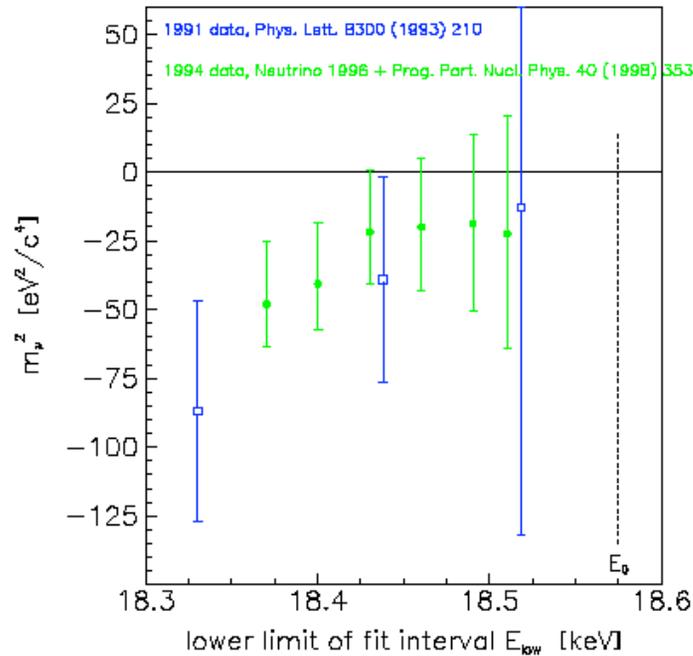
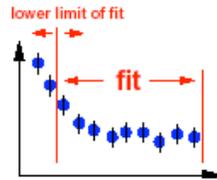
95-97 significant background reduction, signal improvement





Mainz Systematics Resolved

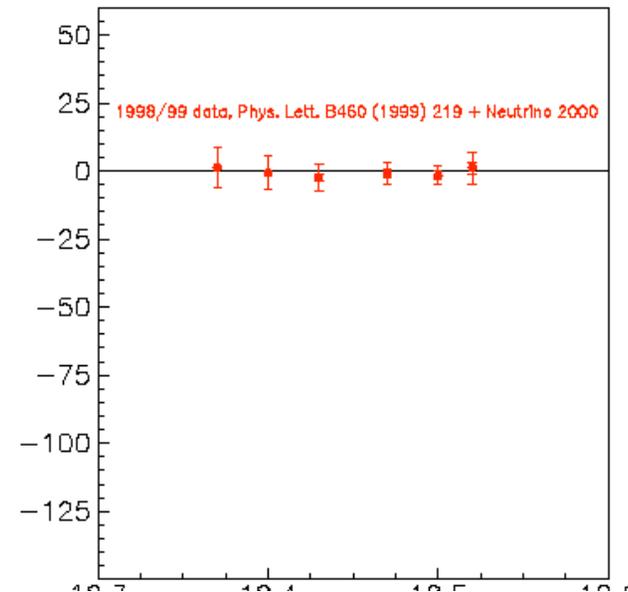
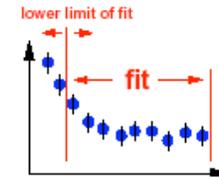
Former problem of negative m_ν^2 at Mainz



⇒ Problem of missing energy loss
was caused by roughening transition

⇒ should be solved by much lower T_2 temperatures

Former problem of negative m_ν^2 at Mainz



$T_{source} = 1.8$ K
Trap pulsing on

Mainz Results



Recent runs (Q5 and greater) exhibit good reduced χ^2 and are stable over a varying fit interval.

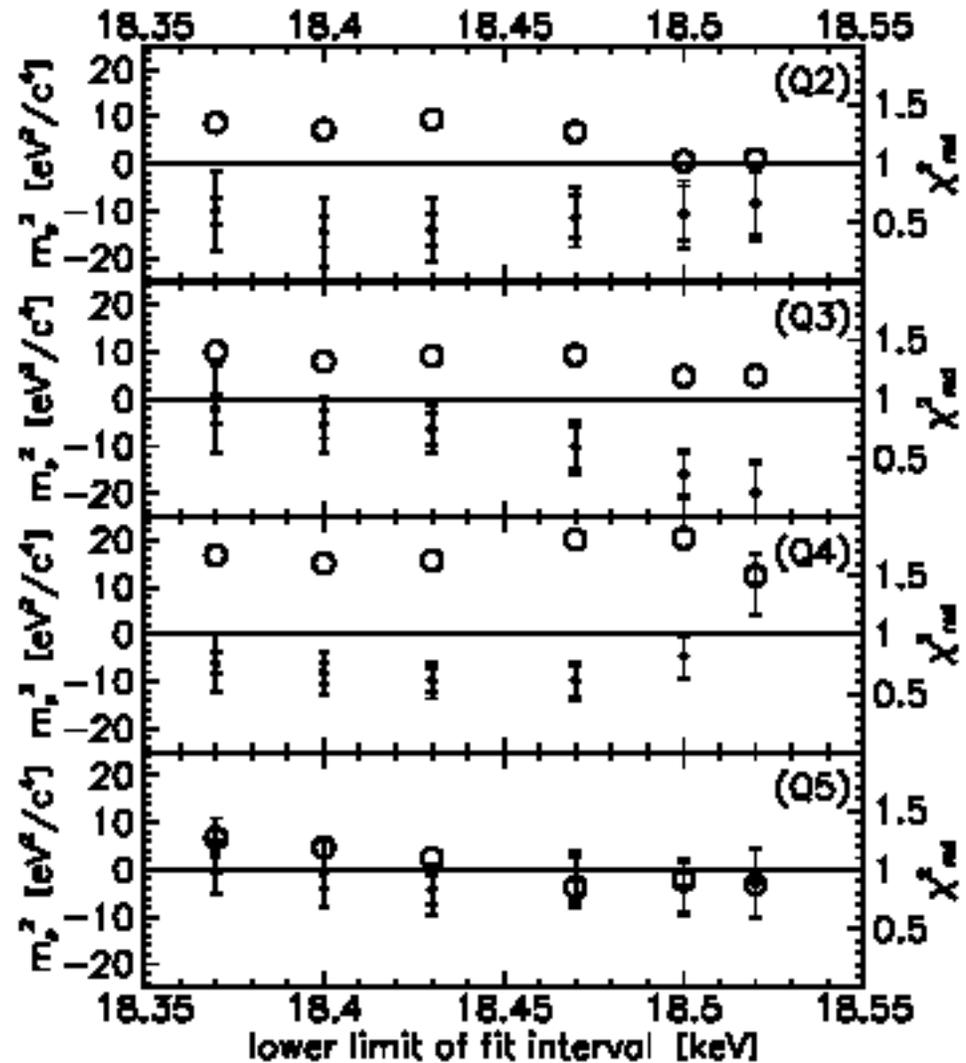
Change made to “sweep” spectrometer backgrounds between data points starting at run Q5.

Detailed studies published on source systematics.

$$m_{\square}^2 = -1.6 \pm 2.5 \pm 2.1 \text{ eV}^2$$

$$m_{\square} \leq 2.2 \text{ eV (95\%CL)}$$

A solid result.



Future β -decay endpoint measurements

- **Ultimate sensitivity of spectrometers**
 - require instrumental resolution of $\sim m_\beta/E_o$
 - spectral fraction per decay that falls in the last m_β of the spectrum is $\sim (m_\beta/E_o)^3$
 - source thickness is set by the inelastic scattering cross-section ($3.4 \times 10^{-18} \text{ cm}^2$), $\Delta n \leq 1$
 - If one wants ~ 1 event/day in last m_β of the spectrum
 - for a 10 m magnetic spectrometer $m_\beta \sim 1.7 \text{ eV}$
 - for a 3 m dia. solenoid retarding field spectrometer $m_\beta \sim 0.3 \text{ eV}$
- **Calorimetric detector sensitivity**
 - evade source-thickness limit, because no e-loss problem
 - limited by response time, and eventually pileup
 - requires fine segmentation, many detectors

See Wilkerson and Robertson, Direct Measurements of Neutrino Mass, Sect 3.6

planning the next-generation direct ν mass experiment

experimental observable in β -decay is m_ν^2

aim : improvement of m_ν by **one** order of magnitude (3 eV \rightarrow 0.3 eV)

requires : improvement of m_ν^2 by **two** orders of magnitude (9 eV² \rightarrow 0.09 eV²)

improve statistics :

- stronger tritium source (factor 40) (& larger analysing plane)
- longer measuring period (~ 100 days \rightarrow ~ 1000 days)

improve energy resolution :

- large electrostatic spectrometer with $\Delta E = 1$ eV (factor 4 improvement)

but : count rate close to β -end point drops very fast ($\sim \delta E^3$)

*last 10 eV : 2×10^{-10}
last 1 eV : 2×10^{-13} of total β -intensity*

Karlsruhe Tritium Neutrino Experiment (KATRIN)

arXiv:hep-ex/0109033

next-generation experiment with *sub-eV* neutrino mass sensitivity

FH Fulda - FZ & U Karlsruhe - U Mainz - INP Prague - INR Troitsk - U Washington

high luminosity

background suppression

high energy resolution

control of systematics

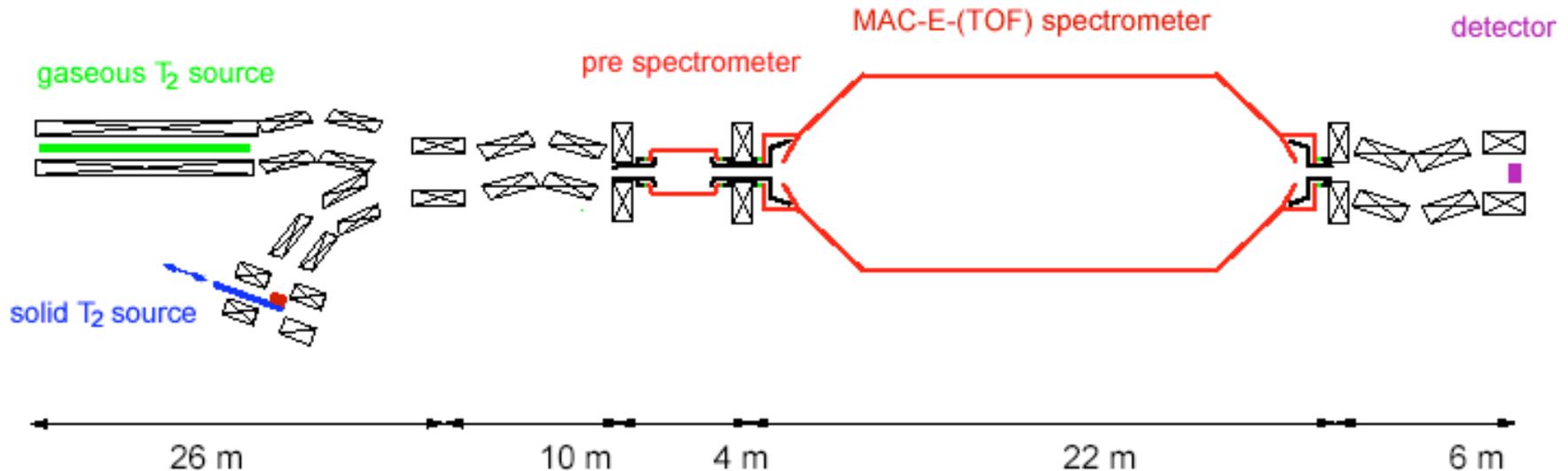
molecular tritium source

pumping

pre-filter

energy analysis

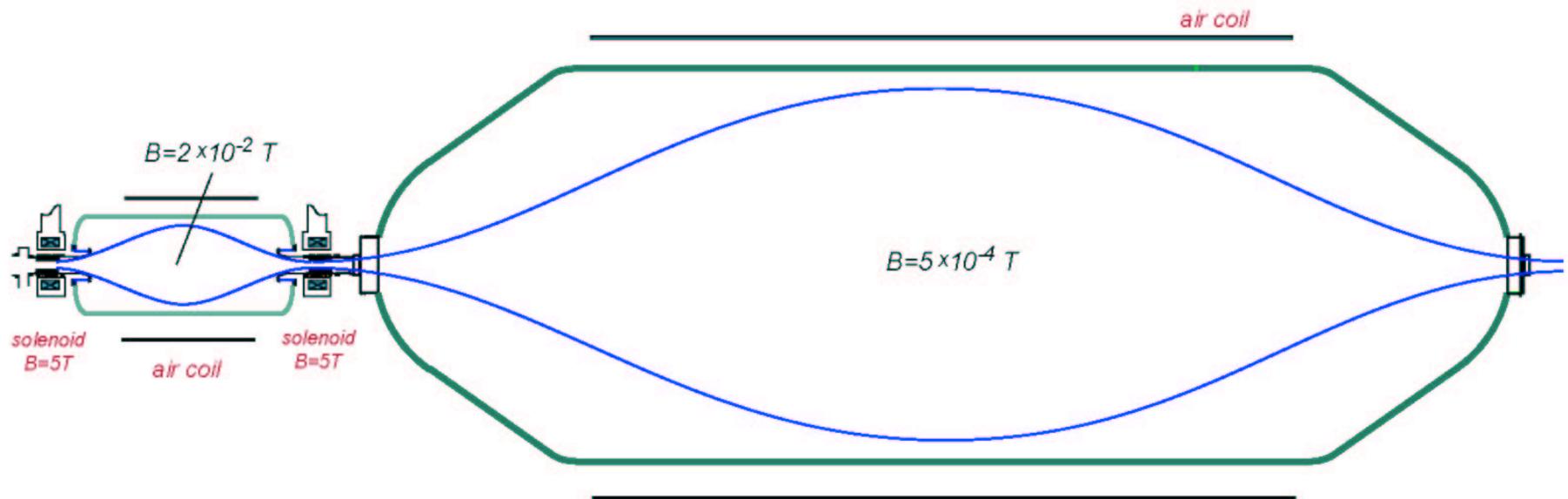
fi-electron counting



electrostatic spectrometers - properties and geometry

electrostatic analysis of tritium β -decay electrons (electrode system)

XUHV - conditions : $p < 10^{-11}$ mbar (degassing rate 10^{-13} mbar l / cm² s)



pre-spectrometer

fixed retarding potential 18.4 kV

$\varnothing = 1.7$ m / $L = 4.0$ m

$\Delta E = 80$ eV

main spectrometer

variable retarding potential 18.5-18.6 kV

$\varnothing = 7$ m / $L = 20$ m

$\Delta E = 1$ eV



Technological Challenges

electrostatic spectrometer

construction large vessel ($\varnothing=7\text{m}$, $l=20\text{m}$)

XHV ($p < 10^{-11}$ mbar)

HV control & stabilization

optimized electrode system

electron transport

> 30 superconducting solenoids

lHe and lN₂ supply (200W cooling power)

optimized particle tracking ($l > 60$ m)

reliable extinction of tritium (freeze out)

tritium sources

stable & safe tritium supply

high luminosity & reliability

control of syst. effects (TOF op., calib.)

solid state detector

excellent $\Delta E/E$ in high B-field ($< 1\text{keV}$)

good position resolution

mK operation of bolometer

experiment will be operational for several years

interdisciplinary solutions are required



Forschungszentrum Karlsruhe



KATRIN & TLK

IK

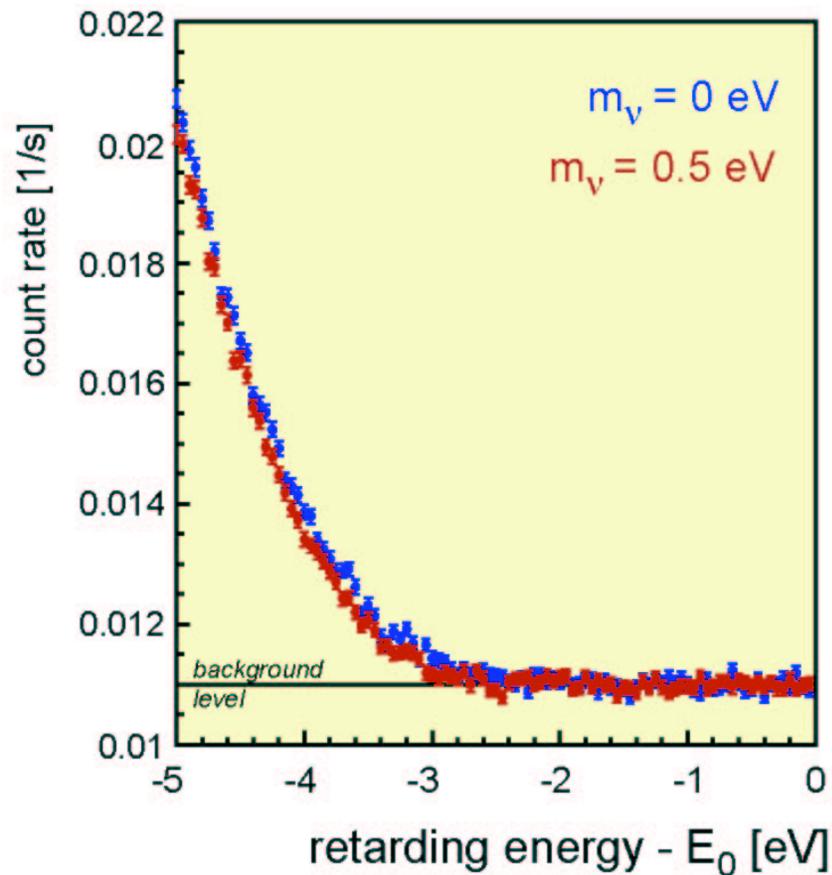
ITP



Estimated KATRIN sensitivity for neutrino masses

realistic MC simulation of sub-eV ν -mass signal close to sensitivity limit

narrow interval close to β end point (last 5 eV) from WGTS



input parameters for simulation :

measuring time : 3 years

$\Delta E = 1 \text{ eV}$ (spectrometer)

background rate = 11 mHz

WGTS :

column density $5 \times 10^{17} / \text{cm}^2$

max. accepted angle 51°

molecular excitations included



Systematic Uncertainties

δE -interval = 15-20 eV

KATRIN focuses on very narrow region below E_0

($\Delta E=1\text{eV}$, high T_2 luminosity): many systematic uncertainties reduced

- **no** contribution from excited electronic states of $^3\text{He-T}$ ($\delta E > 25\text{ eV}$)
- **small** contribution from inelastic scattering
in source (for δE -Interval of 25 eV : 2% of signal from scattered electrons)

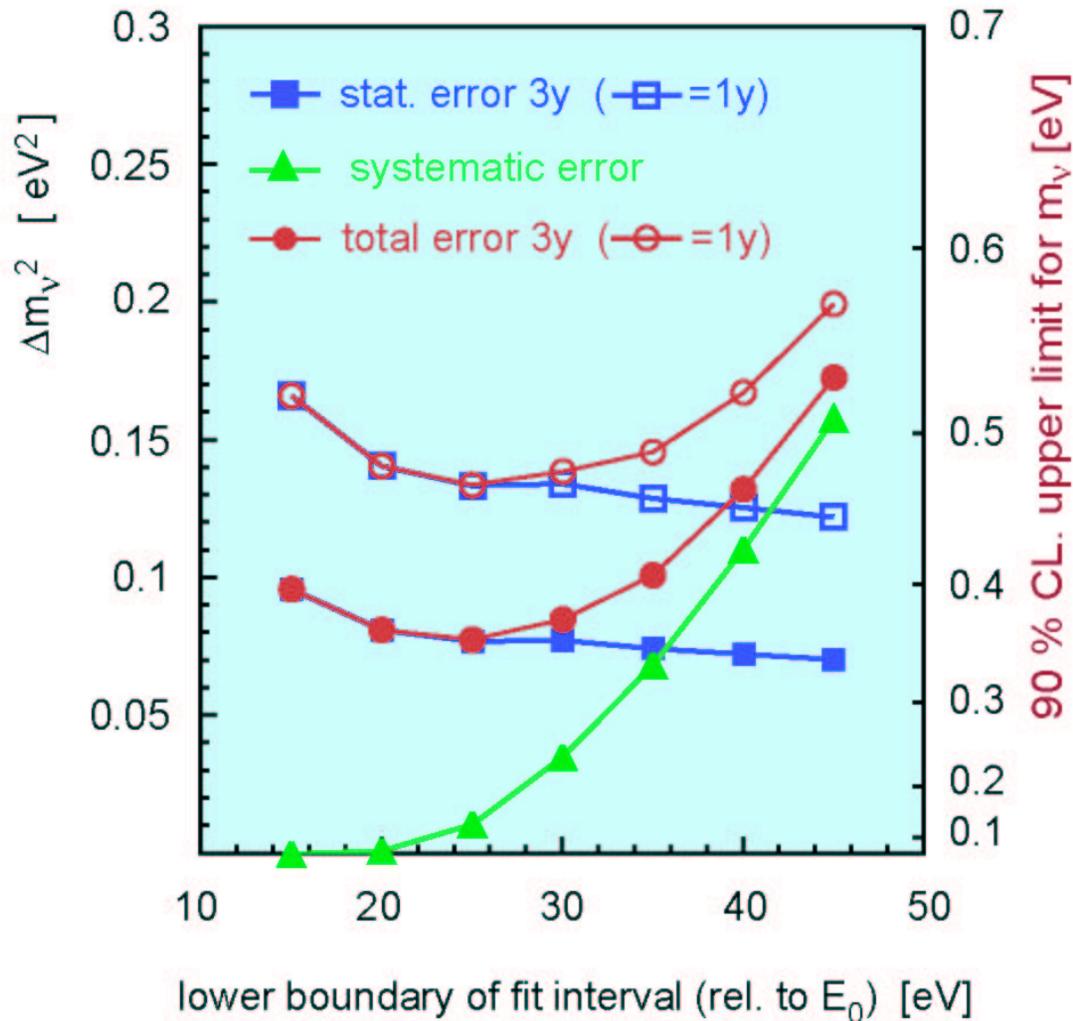
+ better vacuum & higher T_2 purity

remaining uncertainties :

- calculations of rotational-vibrational excitations of $^3\text{He-T}$ ground state (0.2% theory uncertainty)
- inelastic scattering of β -electrons in WGTS
(2% uncertainty on σ_{tot} , can be improved)
- solid state effects (self-charging of film, neighbour excitations, ...) only QCTS
- stability of settings : HV calibration and stabilisation
WGTS activity and T_2 -purity



estimates of KATRIN sensitivity for m_ν



assumptions for simulation:

$\Delta E = 1$ eV (spectrometer)

background rate = 11 mHz

WGTS : $\rho d = 5 \times 10^{17} / \text{cm}^2$

area = 29 cm^2

max. accepted angle 51°

systematic error :

2% energy loss in WGTS

$m_\nu < 0.35$ eV (90% CL.)



KATRIN - time schedule

- 1/2001 first presentation at international workshop at Bad Liebenzell
- 6/2001 formal founding of KATRIN collaboration
- 9/2001 Letter of Interest (LoI) submitted hep-ex/0109033
BMBF funding 'astroparticle physics' for german universities
- 5/2002 Complete successful FZK International Panel Review
- 12/2002 Submission of proposal
- 2002-03 systematic studies of background processes and design optimisation
funding requests (HGF, DOE, ...) and reviews
pre-spectrometer measurements and R&D studies
- 2004-06 set up of spectrometer, solenoid system, transport system, detector
and tritium sources, hall construction, cryo supply
- 2006 commissioning and begin of data taking



Summary

- All 3 ν masses can be probed through the electron neutrino
- Present direct limit on mass of ν_1 , ν_2 , and ν_3
 - 2.2 eV
 - Sum ≤ 6.6 eV
- Direct measurements combined with oscillation and $0\nu\beta\beta$ decay results can discriminate between a variety of 3 and 4 neutrino mass spectrum scenarios.
- KATRIN should be able to achieve sub-eV mass sensitivity of 0.30 eV, nearly an order of magnitude improvement over current direct measurements.
- Sub-eV direct lab limits yield model independent results that will likely be serve as strong constraints on astrophysics.