

Impact of Non-Gaussian errors on CMB lensing and Dark Energy observables in the lensed CMB

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References: Smith, Hu and Kaplinghat, astro-ph/0402442
Smith, Kaplinghat and Hu, to appear

What information is in the primary (unlensed) CMB?

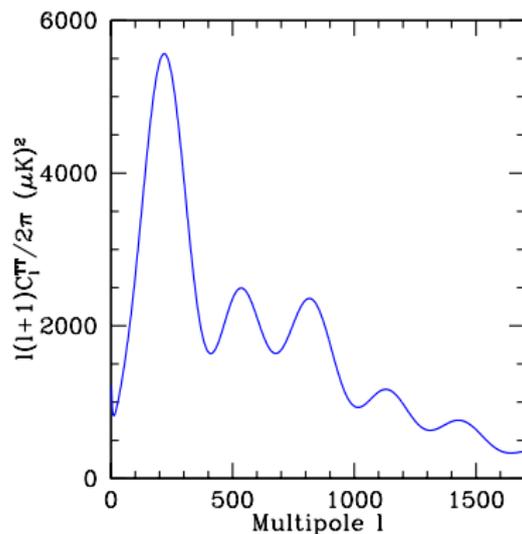
From shape of the power spectra, get strong constraints on:

$$\{\Omega_b h^2, \Omega_m h^2, n_s\}$$

Adding polarization at low multipoles, can separate:

$$\{\tau, \ln(\delta_\zeta)\}$$

(Without this information, $(\ln(\delta_\zeta) - \tau)$ is very well measured, but only this combination.)



What information is in the primary (unlensed) CMB?

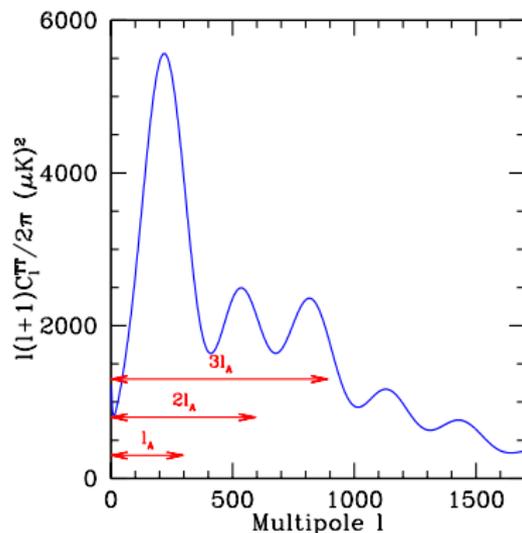
Finally, can measure angular scale of the acoustic peaks:

$$\ell_A = \pi \frac{D_*}{s_*} \quad \leftarrow \text{Angular diameter distance to recombination}$$
$$\quad \quad \quad \leftarrow \text{Sound horizon at recombination}$$

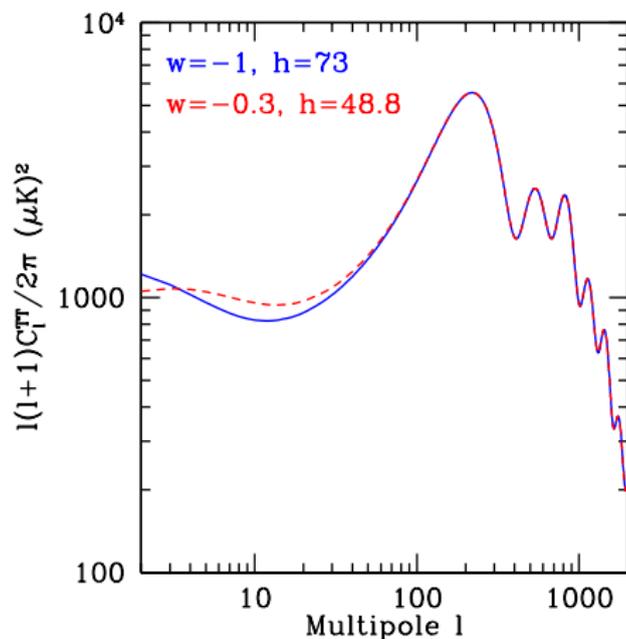
This can be interpreted (via the distance measurement D_*) as one constraint in a dark energy parameter space.

e.g. Hubble parameter h is well measured if $w = -1$ is assumed
(WMAP3: $h = 0.73 \pm 0.03$)
or prior on h constrains w

But, because there is only one observable, large degeneracies exist.



Angular diameter distance degeneracy (in unlensed CMB)



Models with $w = -1$ (and $h = 0.73$) and $w = -0.3$ (and $h = 48.8$) are nearly degenerate in the unlensed CMB

Gravitational lensing

Intervening matter between the surface of last scattering and an observer today lenses the CMB:

$$\tilde{T}(\hat{n}) = T(\hat{n} + \nabla\phi(\hat{n})),$$

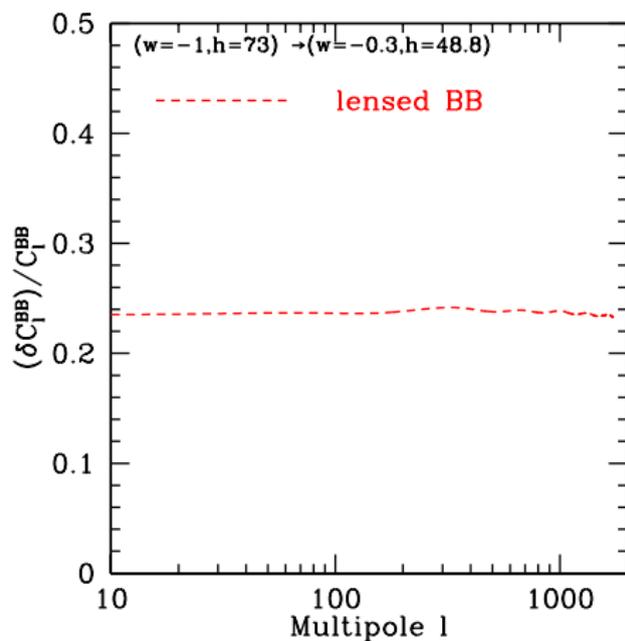
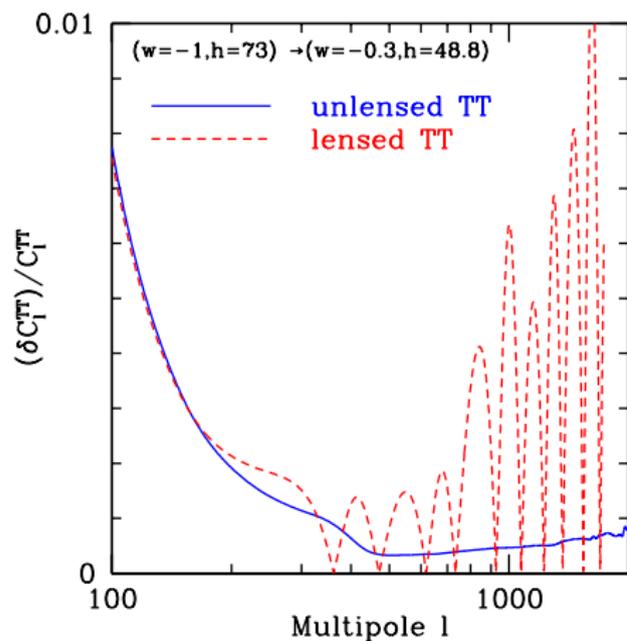
where ϕ is the projected potential, given by the line-of-sight integral

$$\phi(\hat{n}) = -2 \int dD \frac{(D_s - D)}{DD_s} \Psi(D\hat{n}, D)$$

\Rightarrow The CMB has some sensitivity to structure at $z \sim 1$

Lensing breaks the angular diameter distance degeneracy

Fractional difference between $w=-1$ and $w=-0.3$ models:



Scope of talk

Questions:

- ▶ What dark energy observables are contained in the lensed CMB? (complementing $\ell_A = \pi D_*/s_*$)
- ▶ Are parameter constraints from the lensed CMB affected by non-Gaussian statistics?

Caveats:

- ▶ Parameter forecasts computed in the Fisher matrix approximation
- ▶ Only parameter constraints from power spectra (i.e., two-point statistics) are considered. Higher-order statistics (such as lens reconstruction methods) may improve constraints.

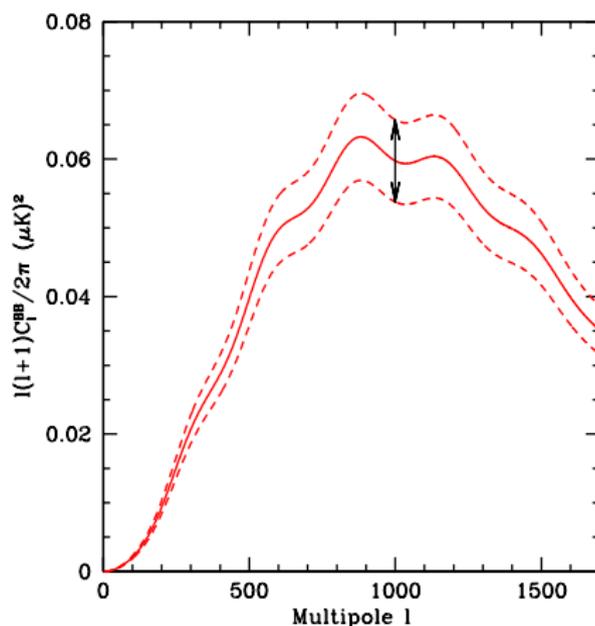
Why is the lensed CMB non-Gaussian?

$$\begin{aligned}\tilde{T}(\hat{n}) &= T(\hat{n} + \nabla\phi(\hat{n})) \\ &= T + (\nabla^a\phi)(\nabla_a T) + \frac{1}{2}(\nabla^a\phi)(\nabla^b\phi)(\nabla_{ab}T) + \dots\end{aligned}$$

- ▶ Terms starting with the second are non-Gaussian (products of Gaussian fields)

Non-Gaussianity: intuitive argument

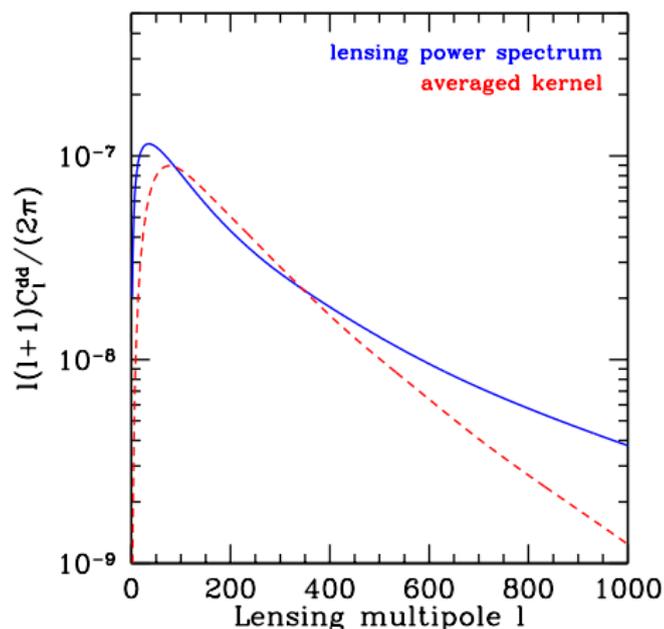
Consider the question: how well can the overall amplitude of the lensing B-modes be measured?



Cosmic variance limited, $l_{max} = 2000$: 0.07%

Non-Gaussianity: intuitive argument

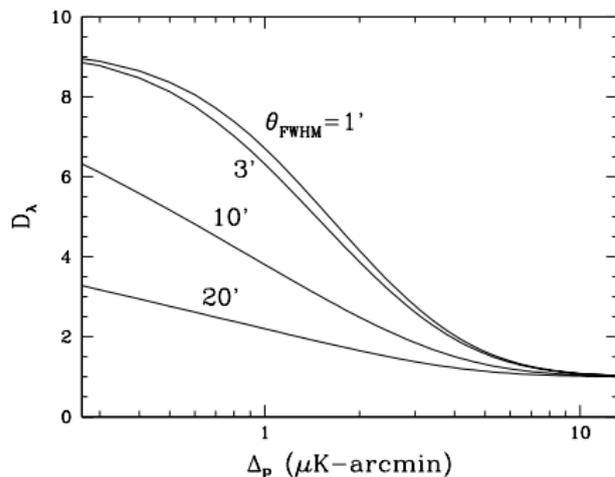
Now ask: Which angular scales in the lensing potential contribute to the overall B-mode amplitude?



Sample variance limit: 0.20% (assuming fictitious direct measurement of ϕ)

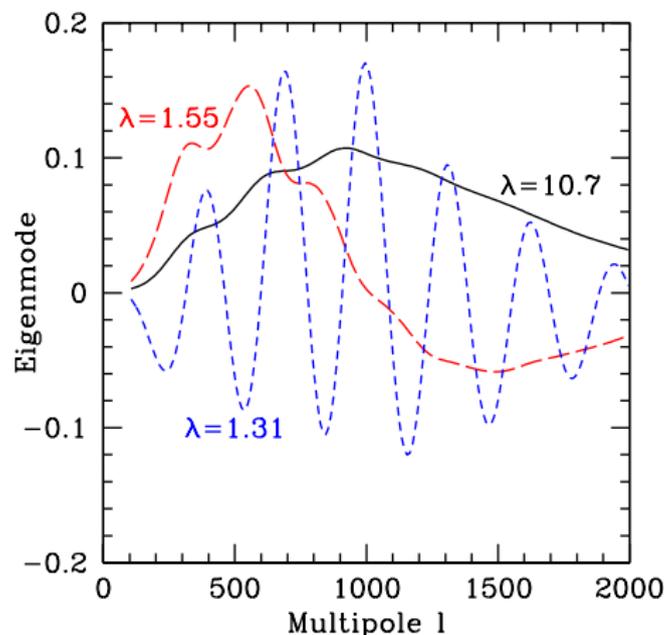
Apparent contradiction

- ▶ Seem to get better constraint from measuring the lensed B-modes (0.07%) than measuring the unlensed CMB *and* the lensing potential ϕ (0.20%).
- ▶ However, the analysis has implicitly assumed Gaussian statistics for BB. When non-Gaussian contributions to the BB power spectrum covariance are included, the variance of the overall amplitude degrades (by the factor D_λ shown).



Non-Gaussianity: complete treatment

- ▶ Non-Gaussianity is always negligible in $\{TT, TE, EE\}$
- ▶ In BB, non-Gaussianity increases the variance of a few eigenmodes



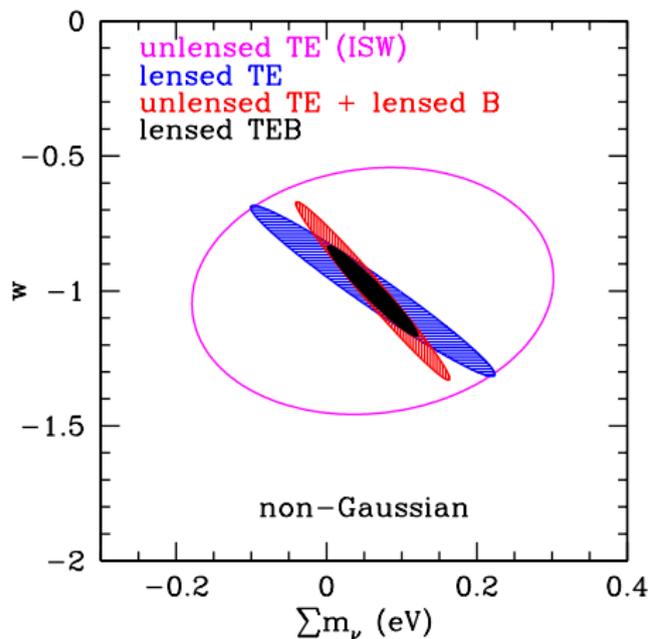
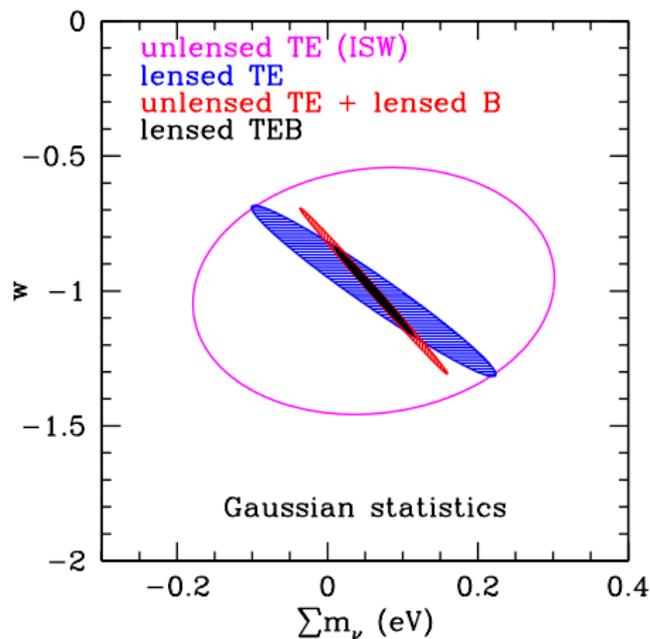
- ▶ Noise level has to be very good for non-Gaussianity to be important!

Parameter Estimation Example (fixed $\Omega_m h^2$, $\ln(\delta_\zeta)$)

Reference survey: 20 $\mu\text{K-arcmin}$, all-sky, zero beam

+ 1 $\mu\text{K-arcmin}$, $f_{\text{sky}} = 0.1$, zero beam

Gaussian vs non-Gaussian uncertainties on ($\Omega_\nu h^2$, w):

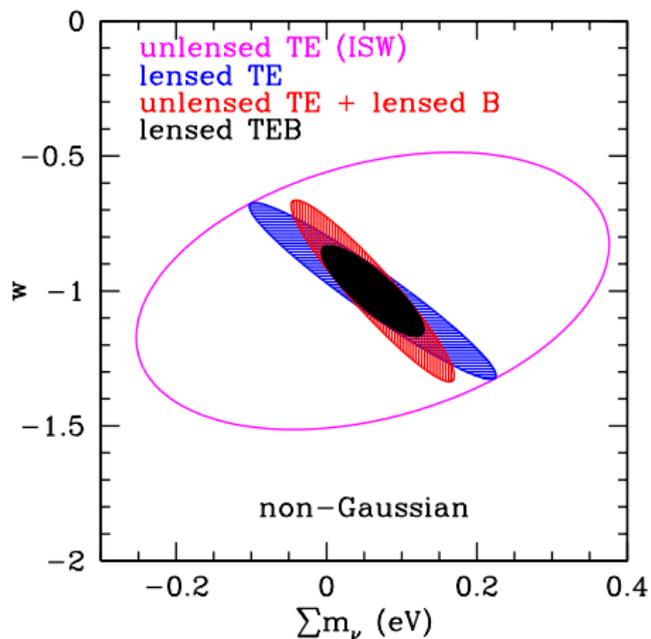
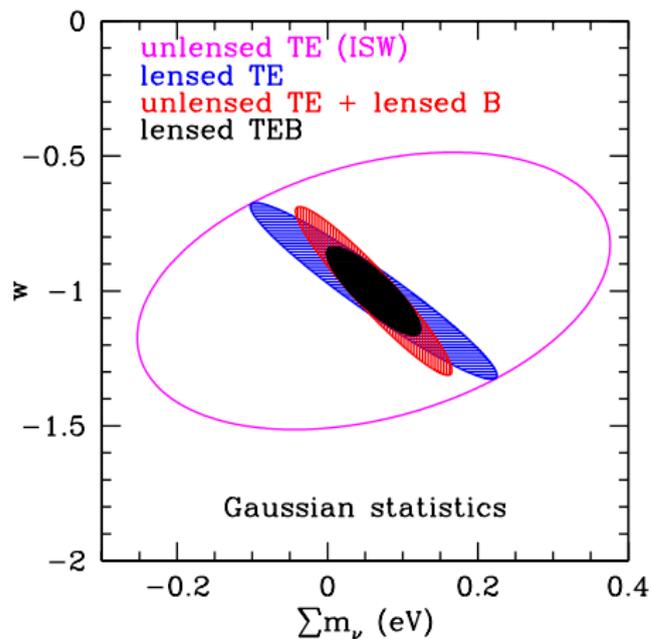


Parameter Estimation Example (fully marginalized)

Reference survey: 20 $\mu\text{K-arcmin}$, all-sky, zero beam

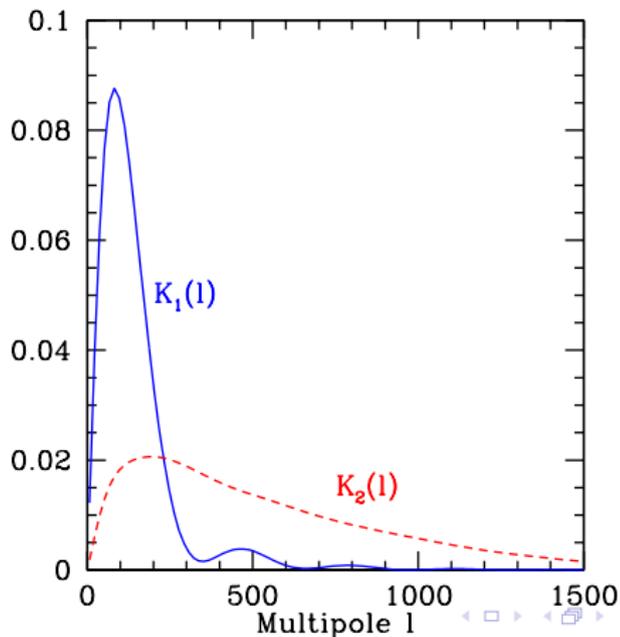
+ 1 $\mu\text{K-arcmin}$, $f_{\text{sky}} = 0.1$, zero beam

Gaussian vs non-Gaussian uncertainties on $(\Omega_\nu h^2, w)$:



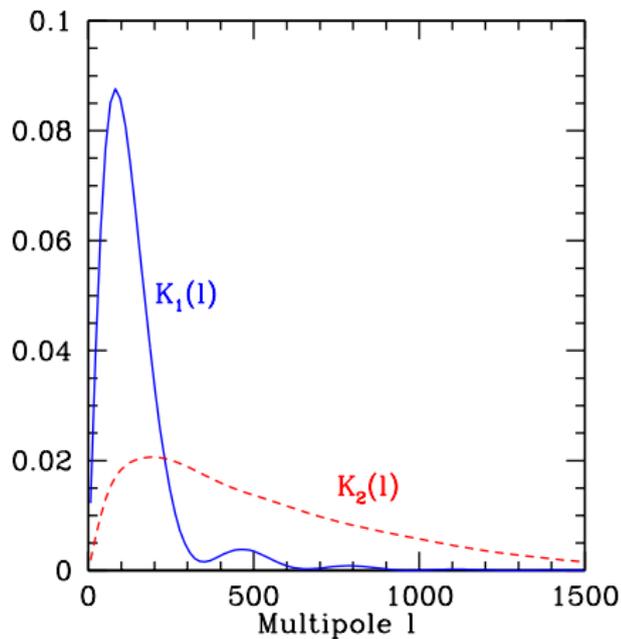
Which modes in $C_{\ell}^{\phi\phi}$ can CMB lensing constrain?

- ▶ Principal component analysis
- ▶ $\{TT, TE, EE\}$ constrain one principal component $K_1(\ell)$ at $\ell \sim 100$
- ▶ BB constrains a distinct principal component $K_2(\ell)$ across a wide range of ℓ

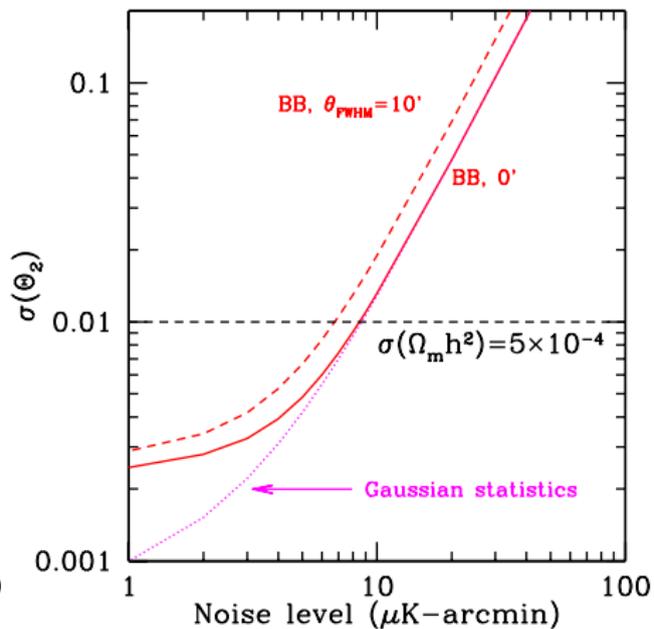
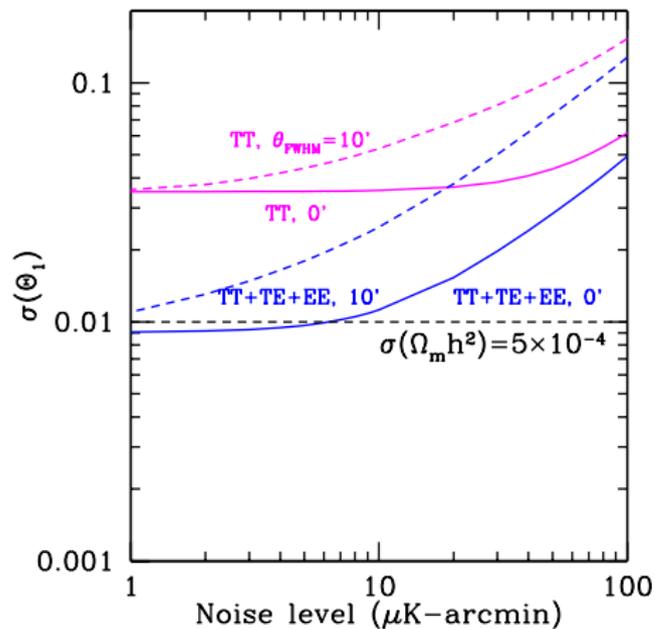


Dark energy observables defined

$$\Theta_1 = \sum_{\ell} C_{\ell}^{\phi\phi} K_1(\ell) \quad \Theta_2 = \sum_{\ell} C_{\ell}^{\phi\phi} K_2(\ell)$$



Dark energy observables: sensitivity

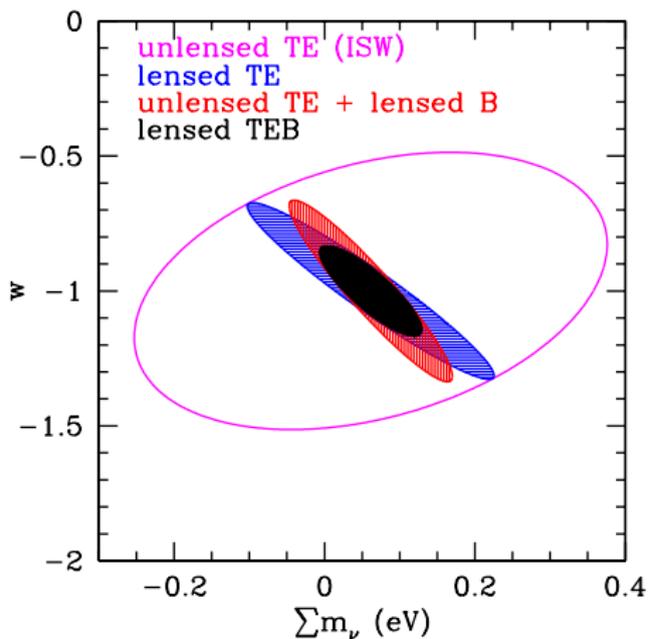
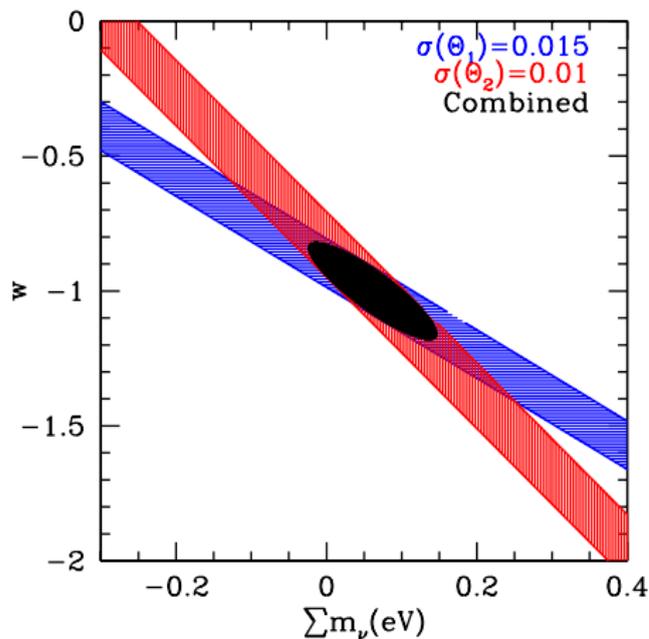


Dark energy observables: example

Reference survey: 20 $\mu\text{K-arcmin}$, all-sky, zero beam

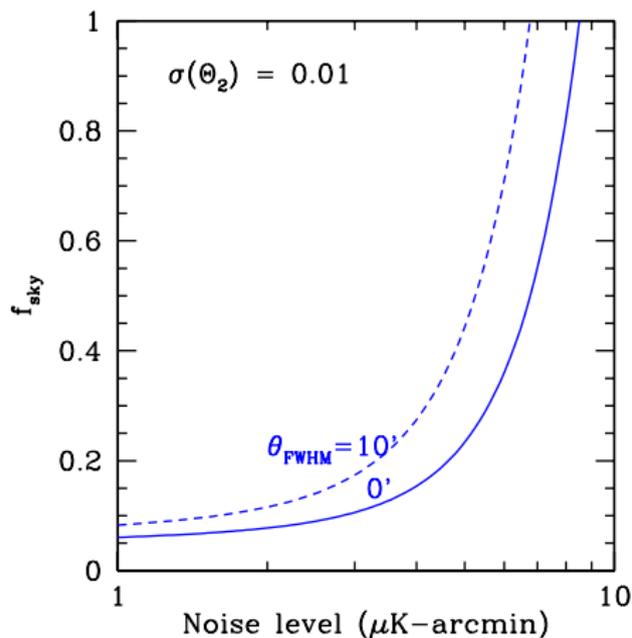
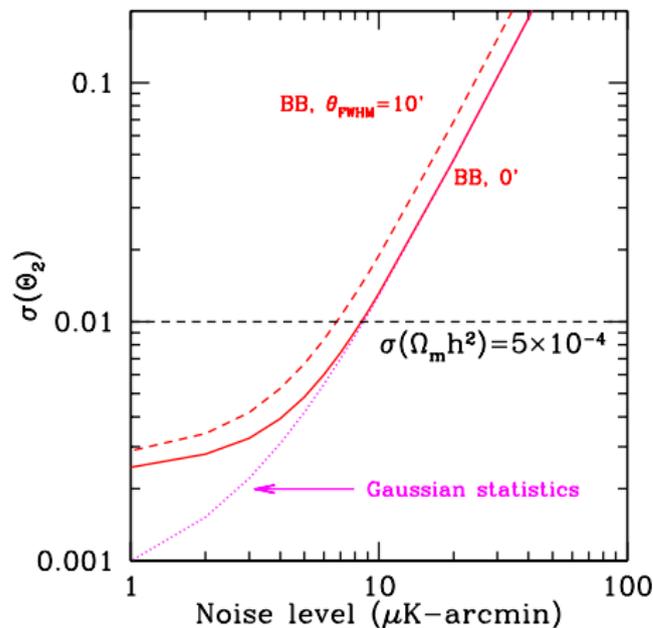
+ 1 $\mu\text{K-arcmin}$, $f_{\text{sky}} = 0.1$, zero beam

$\{\Theta_i\}$ picture vs complete Fisher calculation:



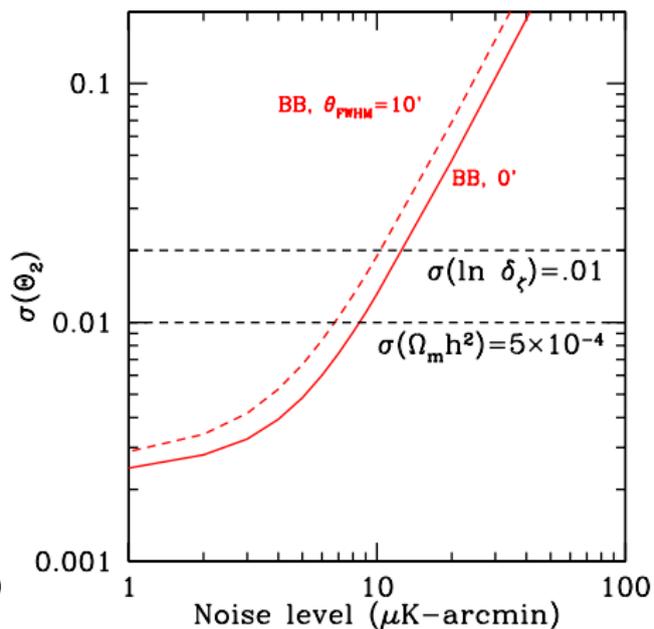
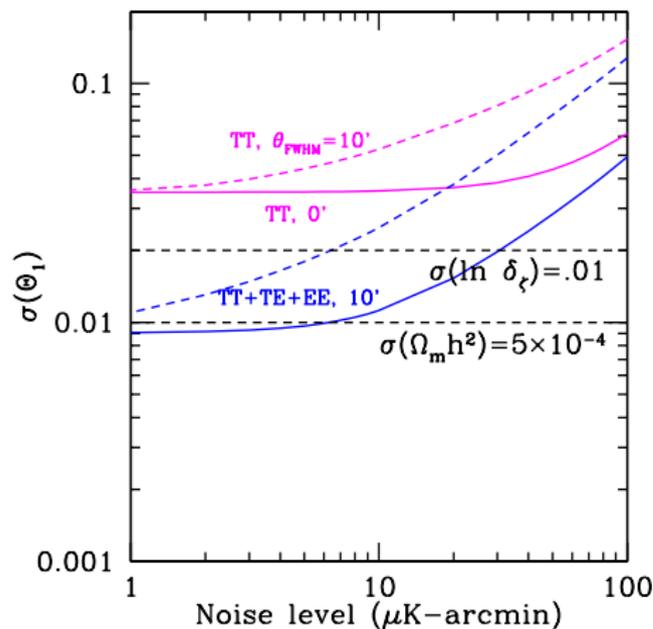
Application 1: existence of a B-mode sensitivity “floor”

- ▶ Uncertainty in $\{\Omega_m h^2\}$ makes a ~ 0.01 contribution to $\sigma(\Theta_2)$
- ▶ Improving experimental sensitivity beyond this level will not improve dark energy constraints



Application 2: reionization history uncertainties

- ▶ Without assuming sharp reionization, $\sigma(\tau) \sim 0.01$ is best possible
- ▶ Uncertainties in Θ_i from $\ln(\delta_\zeta)$ (rather than $\Omega_m h^2$) would then be the limiting factor at high sensitivity



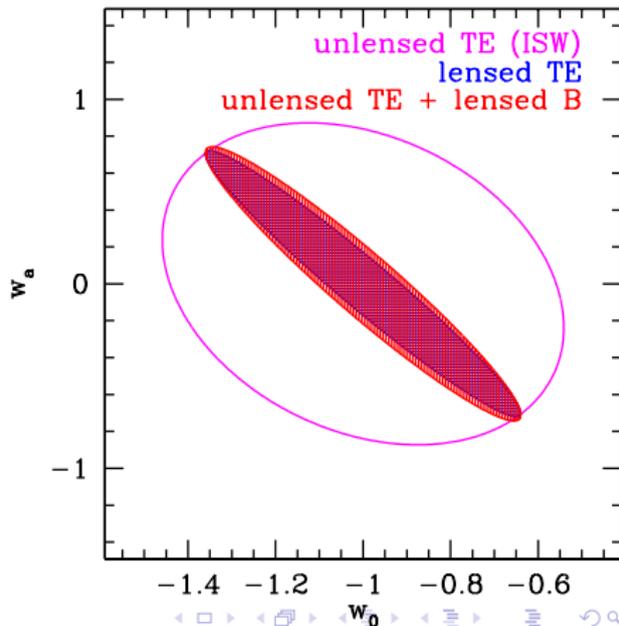
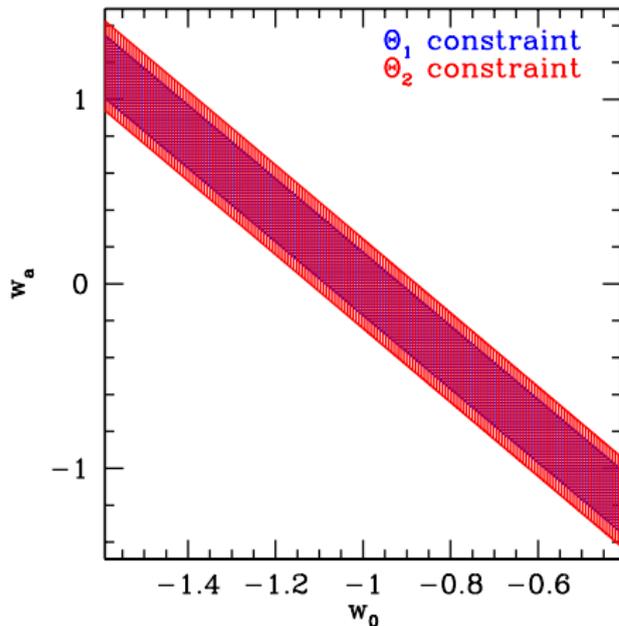
Application 3: constraints on (w_0, w_a)

Fixing neutrino mass, but allowing $w(a) = w_0 + (1 - a)w_a$.

Reference survey: 20 $\mu\text{K-arcmin}$, all-sky, zero beam

+ 1 $\mu\text{K-arcmin}$, $f_{\text{sky}} = 0.1$, zero beam

$\{\Theta_i\}$ picture vs complete Fisher calculation:



Summary

- ▶ The unlensed CMB places excellent constraints on one dark energy observable $\ell_A = \pi D_*/s_*$ but is otherwise degenerate.
- ▶ Through lensing, the CMB is sensitive to density fluctuations at $z \sim 1$, which break the degeneracy.
- ▶ Lensed $\{T,E\}$ can constrain a second observable Θ_1 ; lensed B can constrain a third observable Θ_2 .
- ▶ Non-Gaussianity in lensed $\{TT,TE,EE\}$ is always negligible.
- ▶ Non-Gaussianity in lensed BB degrades the overall amplitude uncertainty by a factor of ~ 10 ; this does not affect dark energy uncertainties after marginalizing $\Omega_m h^2$.
- ▶ There is a B-mode sensitivity “floor”, beyond which dark energy constraints do not improve.
- ▶ The parameters w and w_a cannot be separated using the CMB alone.